

Electroweak Elastic Scattering In Heaven and Earth

Irvine, Dec. 2002

C. J. Horowitz

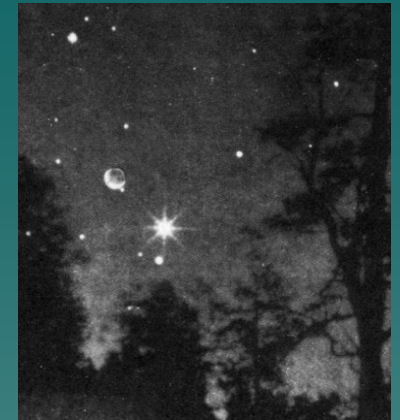
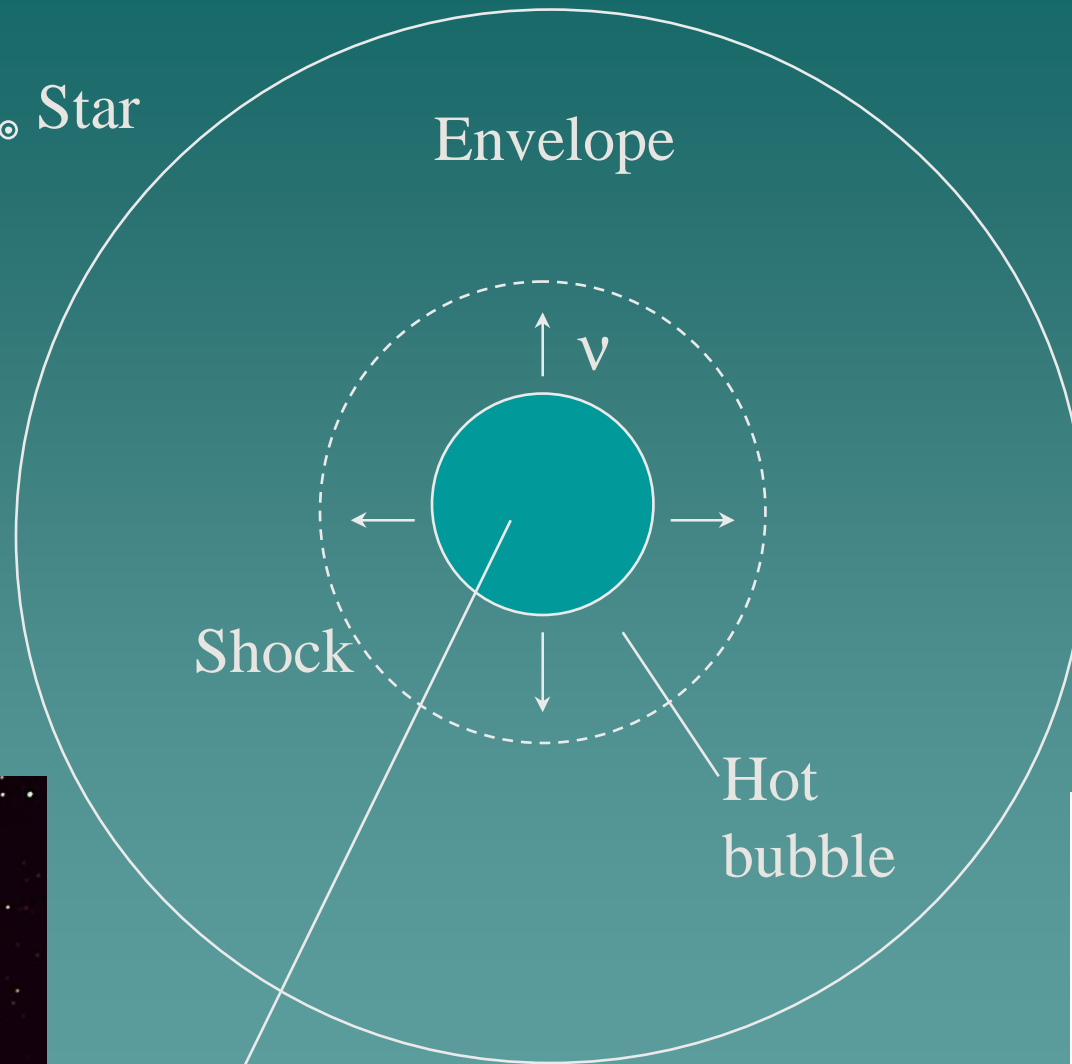
A stylized, teal-colored silhouette of a mountain range is located in the bottom right corner of the slide. The mountains are rendered in a layered, blocky style, with varying heights and peaks, creating a sense of depth and texture against the dark teal background.

Electroweak Elastic Scattering

- ◆ ν interactions in supernovae
 - Opacity dominated by ν -n elastic.
 - ν energies and r-process nucleosynthesis.
 - SN detection via ν -p and ν -A elastic.
- ◆ Strange quark content of nucleon
 - Parity violating electron scattering
 - ν -p elastic scattering to determine strange quark contributions to nucleon spin Δs .
 - Measure ratio of neutral to charged current.

Core collapse supernova

$>8M_{\odot}$ Star

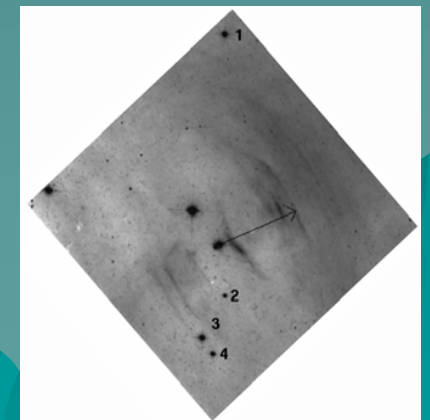


July 5,
1054

Crab nebula



Proto-neutron star: hot, e rich

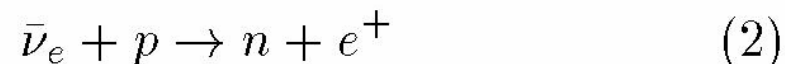


Opacity Dominated by ν -n Elastic

- ◆ Energy transport mostly by ν_x ($\equiv \nu_\mu, \nu_\tau$) because twice as many as ν_e and ν_x without charged currents, have longer mean free path.
- ◆ Opacity of ν_x mostly from ν -n elastic.
- ◆ Incorrect elastic cross sec caused Oak Ridge simulation to explode when it did not with correct one.
- ◆ Uncertainty in ν -n cross sec from strange quarks relevant for SN simulations.

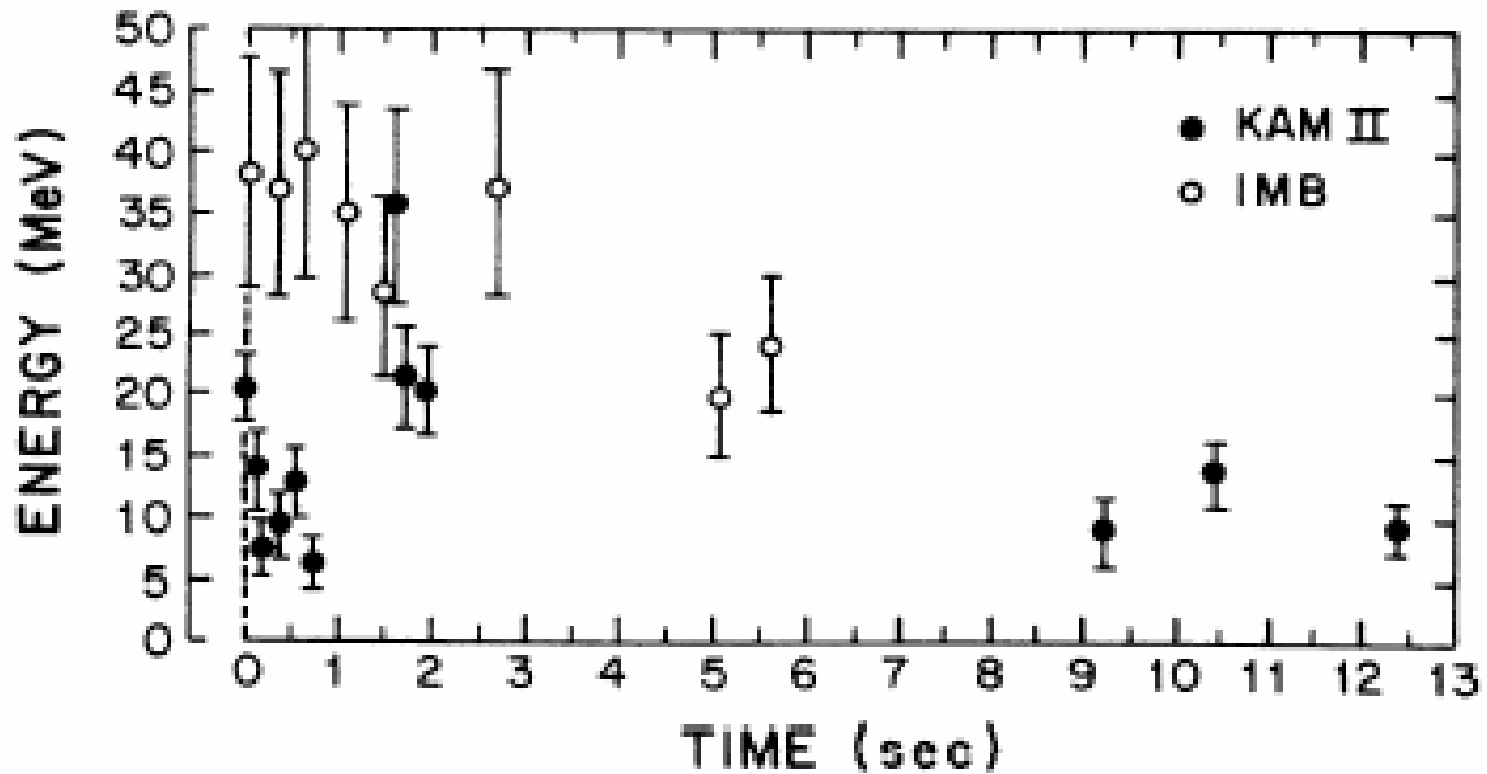
R-process in neutrino driven wind

- ◆ Half of heavy elements made by rapid neutron capture. Best site is ν driven wind in SN.
- ◆ Low density region above proto-neutron star dominated by large ν flux.
- ◆ Initial neutron to proton ratio and Y_e in wind set by relative rates:



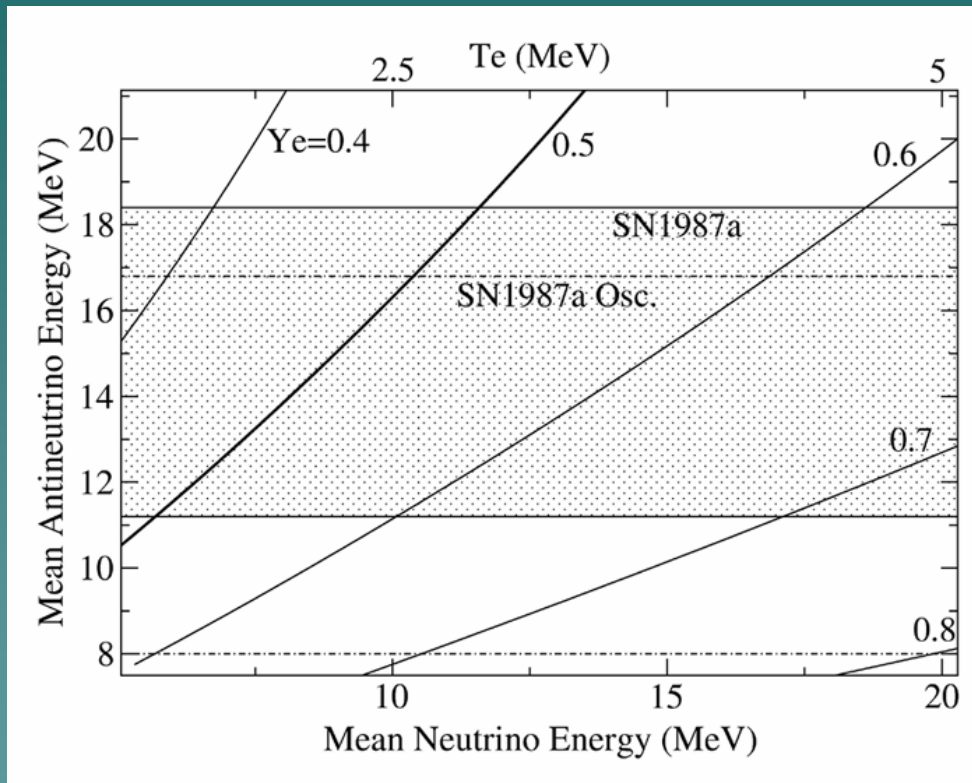
- ◆ Cross section for (1) larger than (2) because of n-p mass difference (Q value) and because of weak magnetism. Note, C violating nature of weak interactions insures ν -nucleon cross sections are larger than anti- ν -nucleon.

Historic IMB SN1987A Data



Twenty $\bar{\nu}_e$ detected via $\bar{\nu}_e + p \rightarrow n + e^+$.

n/p ratio in ν -driven wind



For wind to be neutron rich must be above dark $Y_e=0.5$ line and below solid SN1987A limit line.

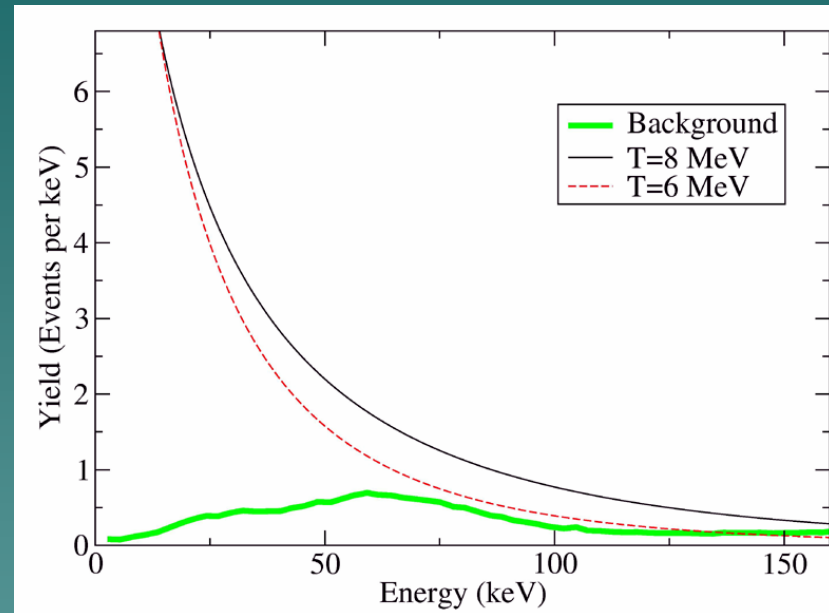
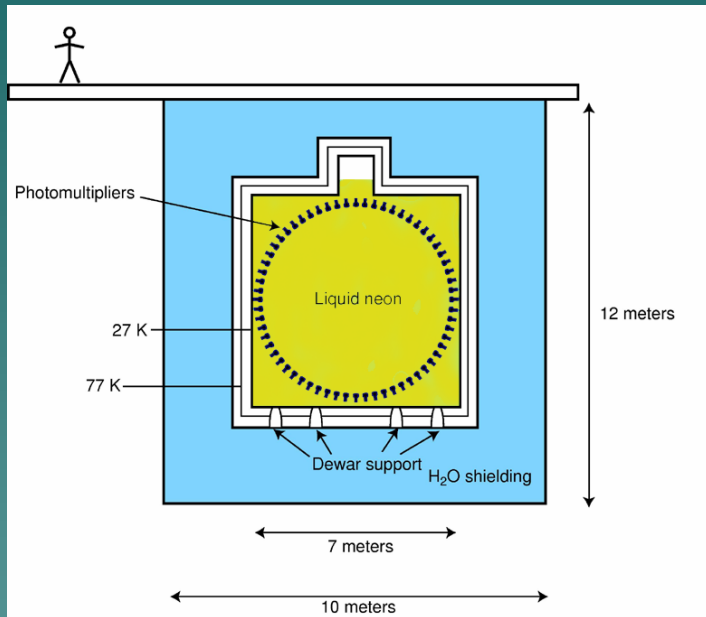
This requires ν_e to be much colder than present simulations.

Present simulations are not robustly neutron rich. Where do n come from for r-process?

Supernova Spectra and ν Oscillations

- ◆ Expect E of $\nu_e < \bar{\nu}_e < \nu_x$ because ν_x have no charged current and star is neutron rich.
- ◆ Difference in E of $\nu_x - \nu_e$ is lever arm for ν osc.
- ◆ Active ν osc don't help r-process. Raising E of ν_e makes problem worse. E of $\bar{\nu}_e$ constrained by SN1987A data.
- ◆ Super-K sees E of $\bar{\nu}_e$ well. Need E of ν_x
- ◆ Elastic ν -p or ν -A scattering can measure energy of ν_x via recoil spectrum.
- ◆ ν -p elastic perhaps in Kamland [J. Beacom]
- ◆ ν -A has very large coherent cross section. Yield of a few or more events per *ton* for a Galactic SN.

Supernova at 10 kpc in CLEAN



CLEAN is detector for low E solar ν via ν -e elastic. Monte Carlo for backgrounds, mostly radioactivity in phototubes. -- Dan McKinsey, Kevin Coakley

Green curve is total low energy background in 10 seconds. Black curve is Galactic SN signal of order 400 ν_x events. Can determine ν_x temperature.

Strange quark content of nucleon

- ◆ Three form factors

- ◆ F_1^s, F_2^s, G_a^s

- ◆ Low Q limits:

- $F_1^s(0)=0, dF_1^s/dQ^2 \rightarrow$ strangeness radius

- $\rho_s,$

- $F_1^s=(\rho_s+\mu_s) Q^2/4M^2$ for small Q^2

- $F_2^s(0)=\mu_s$ strange magnetic moment,

- $G_a^s(0)=\Delta s,$ fraction of nucleon spin carried by s

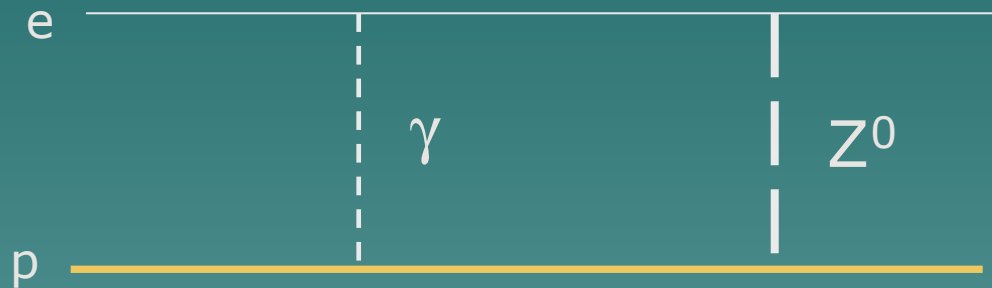
$$\langle f | J_\mu^H | i \rangle = \bar{u}_f(p') \left[\gamma_\mu F_1(Q^2) + i \frac{\sigma^{\mu\nu} q^\nu F_2(Q^2)}{2M_p} + \gamma_\mu \gamma_5 G_A(Q^2) \right] u_i(p),$$

Parity Violating Electron Scattering

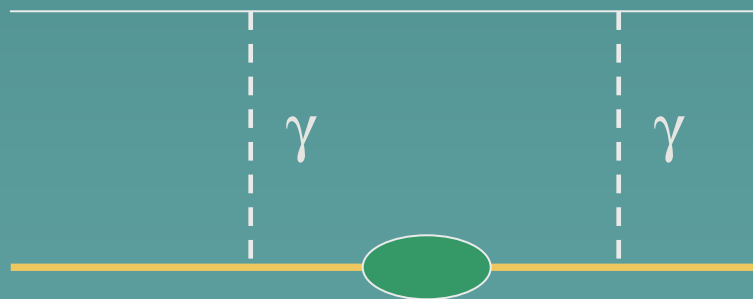
- ◆ Measure asymmetry in cross section of right handed versus left handed electrons.
 $A \approx 1 \text{ ppm}$
- ◆ Sensitive to interference between Z^0 and γ exchange.
- ◆ Forward angles: F_1^s , back angles F_2^s
- ◆ Little sensitivity to G_a^s because of small vec. weak coupling of electron (and radiative corrections). **PV can't get Δs**

Radiative Corrections

- ◆ Lowest order matrix elem. squared



- ◆ Radiative correction

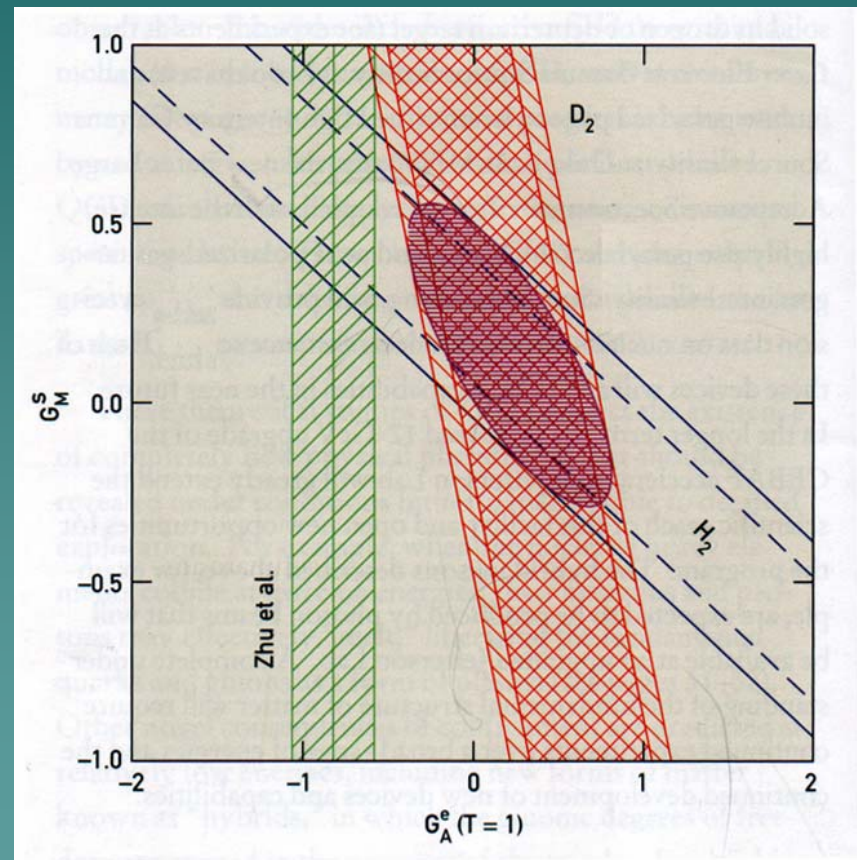


Some Hadronic PV

Example: photon couples to pion loop which has parity violating weak coupling to nucleon

Sample Results

- ◆ PV on H and D give both radiative corrections and μ_S .
- ◆ G_M^S (or μ_S) consistent with zero but large errors.
- ◆ Radiative corrections larger than expected.



Other PV Experiments

- ◆ **HAPPEX**: PV from H at forward angle, $Q^2=0.5$ GeV². $G_E^s+0.4G_M^s=0.02\pm0.03\pm0.02\pm0.03$
- ◆ **G₀** will measure PV from H and D at front and back angles over large Q^2 range. Also PV in **Delta** production.
- ◆ **A4** in Germany, H at intermediate angles.
- ◆ **HAPPEX II**: PV from H at more forward angle $Q^2=0.1$ GeV².
- ◆ Low Q^2 **⁴He** experiment isolates F_1^s
- ◆ Elastic PV from **Pb** will measure neutron radius. [Weak charge of n \gg p].
- ◆ **Q_{weak}** PV from H at low Q^2 as standard model test. Also PV e-e exp. at SLAC.

Bring Me the Head of a Strange Quark

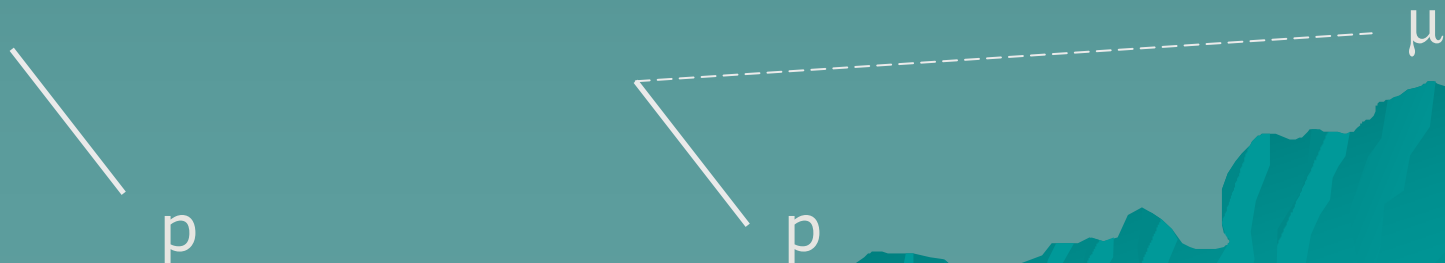
- ◆ Can one find a significantly nonzero strange quark signal with PV electron scattering?
- ◆ It may be very hard to improve significantly on the systematic errors of SAMPLE and G_0 .

Neutrino-Nucleon Elastic Scattering

- ◆ Measure strange quark contributions to nucleon spin Δs
- ◆ Radiative corrections small.
- ◆ Also can measure with antineutrinos.
- ◆ Attractive neutrino fluxes at oscillation beam lines.
- ◆ Need to control systematic errors.

Measure Δs via Ratio of Neutral to Charged Current Scattering

- ◆ Ratio of protons from: $\nu + p \rightarrow \nu + p$ to protons from: $\nu + n \rightarrow \mu^- + p$.
- ◆ Note, both are quasielastic scattering from an $N=Z$ nucleus such as ^{12}C .
- ◆ Very simple observable: ratio of protons of a given E without muons to those with muons.



Example: $E_\nu=0.8$ GeV, $Q^2=0.5$

◆ Neutral to CC ratio $R \approx 0.14$

◆ Error in extracted Δs

– 5% measurement of R	0.04
– ± 0.03 GeV uncer. in M_A	0.01
– ± 0.3 uncer. in μ_s	0.07
– ± 2 uncer. in ρ_s	0.002

[Assume $G_a^s = \Delta s / (1 + Q^2/M_A^2)^2$]

◆ 5% ratio sensitive to Δs at ± 0.04

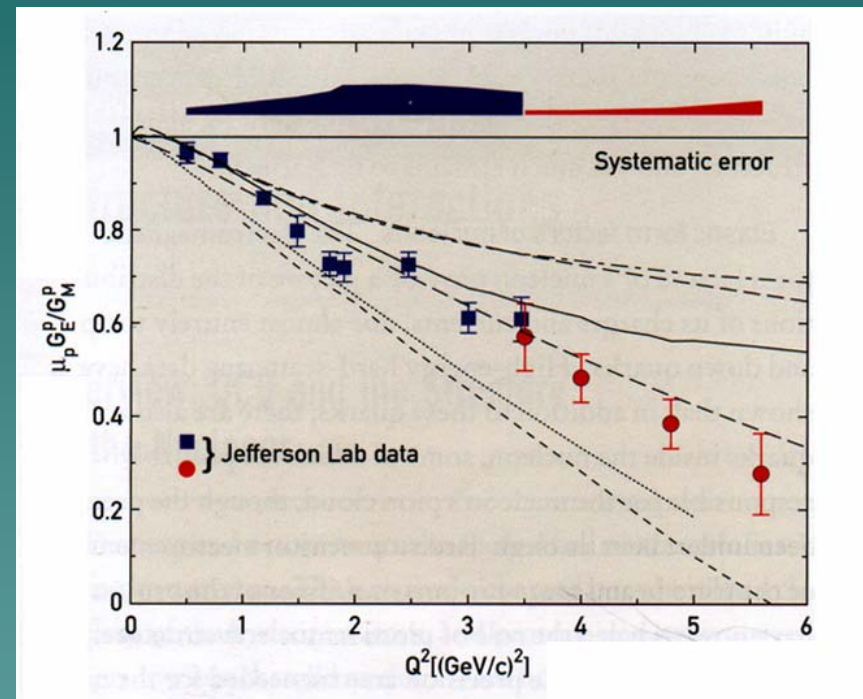
◆ Determine one combination of Δs and μ_s from ν and another from anti- ν .

Experimental Considerations

- ◆ Many systematics, such as absolute flux, proton efficiency, cancel in ratio.
- ◆ Need to identify pions to separate elastic from inelastic events. [This may require a segmented detector.]
- ◆ Possible backgrounds from neutrons and multiple nucleon knockout.
- ◆ Many nuclear structure issues also cancel in ratio, however don't go too low in Q^2 . Want proton recoil energy to be large compared to giant resonances $T_{\text{lab}} > 50-100$ MeV (can use more calculations).
- ◆ Note, more counts and closer to $Q^2=0$ limit for Δs at low Q . Tradeoff in Q^2 choice.

Measurement of Charged Current Axial Form Factor, $g_A/(1+Q^2/M_A^2)^2$

- ◆ Needed to control systematic errors for extraction of Δs .
- ◆ Accurate determination of M_A possible with charged current quasielastic events.
- Search for non dipole behavior? Look at large Q^2 range: $0.5-2+\text{GeV}^2$. How to control systematic errors??



FINESE: Fermi Lab Intense ν Scattering Experiment

- ◆ Proposed near detector on MiniBooNE beam line.
- ◆ Measure Δs via ratio of neutral to charged current.
- ◆ See Rex Tayloe
- ◆ Also Jorge Morfin's proposal for NUMI

