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# DARK MATTER AND DARK SECTORS

Theory Frontier Session, Snowmass 2022

*23 July 2022*

Jonathan Feng, UC Irvine



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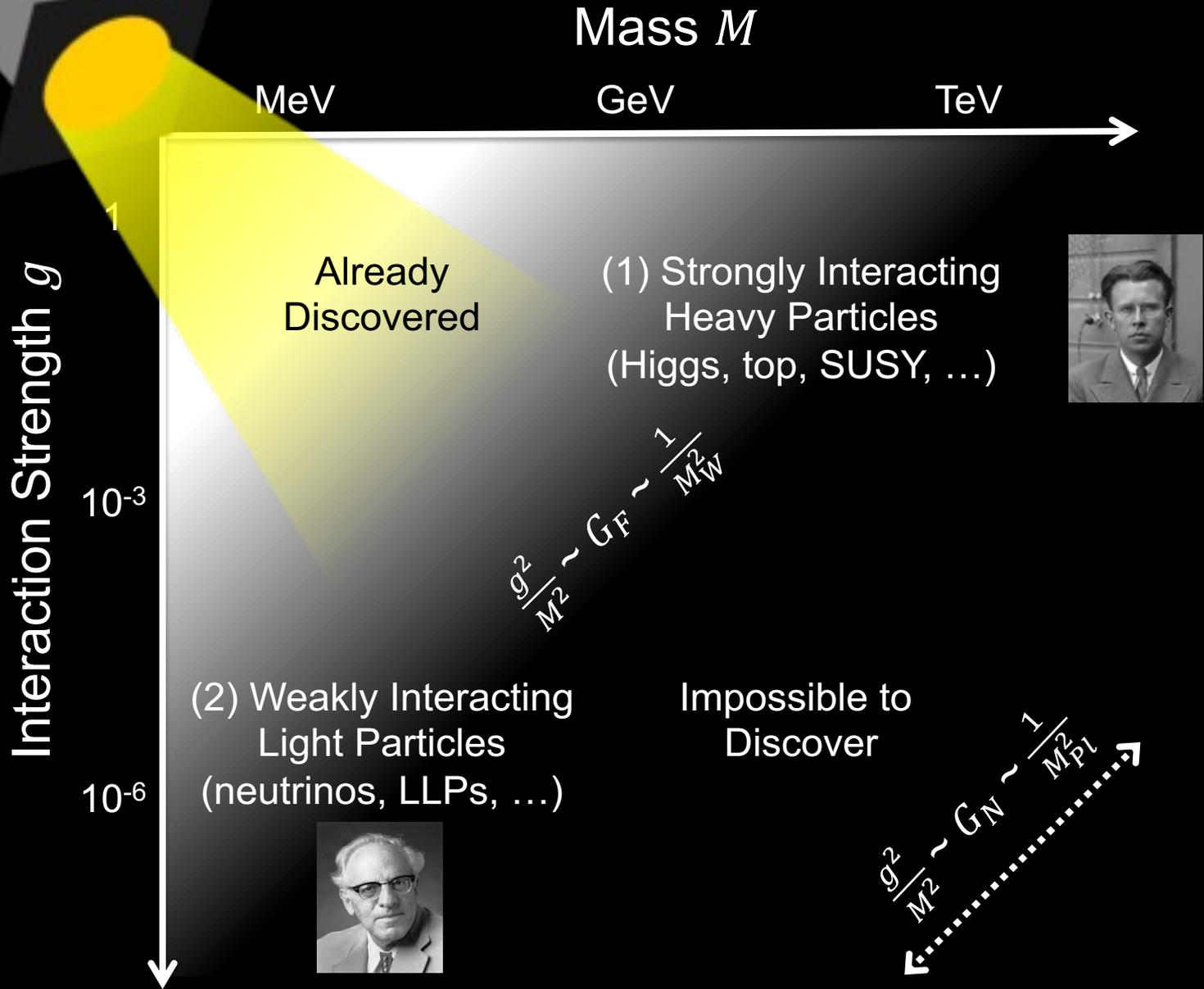


# INTRODUCTION

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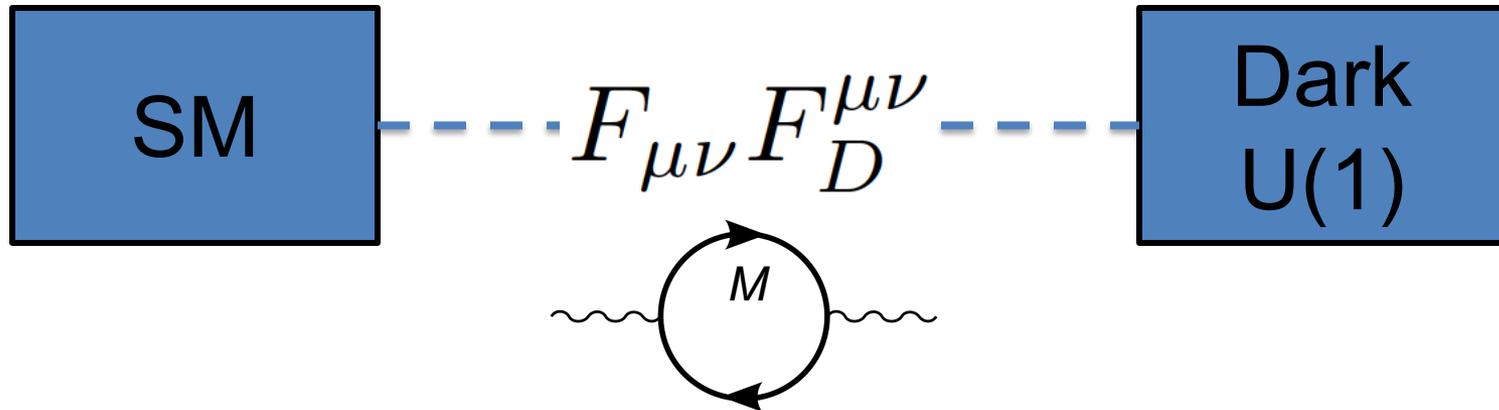
- Dark Matter and Dark Sectors is currently a topic of great interest.
- Many talks at Snowmass, including several talks with almost the identical title:
  - Dark Matter and Dark Sectors: from Theory to Discovery in the Lab, Wednesday, Masha Baryakhtar
  - Dark Matter and Dark Sectors: from Theory to Discovery in the Sky, Wednesday, Joshua Foster
  - Understanding Dark Matter and Dark Sectors at RP Frontier, Wednesday, Stefania Gori
  - Dark Matter and Dark Sectors, Saturday, Jonathan Feng
- Some of the interest follows null results from LHC and ton-scale WIMP searches. But, as I will discuss, there are strong independent reasons to consider dark sectors, and much of the initial interest preceded these null results.
- Connections to almost every frontier, and also to adjacent fields (nuclear physics, condensed matter physics, AMO, and, of course, astrophysics).
- A beautiful example of what can result from the serious and persistent interaction of theorists and experimentalists.

# WHERE IS THE NEW PHYSICS?



# DARK SECTORS

- DM is one of the strongest BSM motivations. In general, it is part of a dark sector. What are its most likely non-gravitational interactions?
- Suppose the dark sector has U(1) electromagnetism. There are infinitely many possible SM-dark sector interactions, but one is induced by arbitrarily heavy mediators:

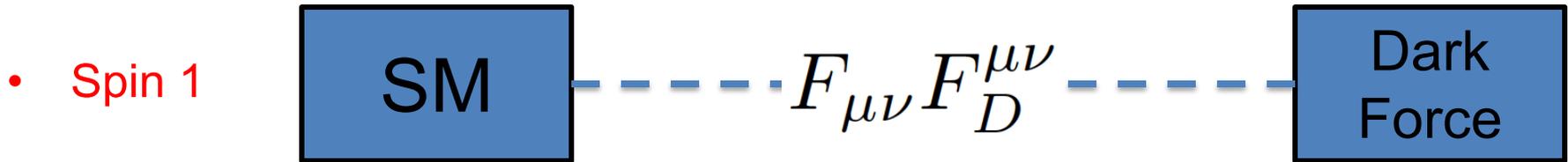


- It is “most likely” because it is non-decoupling. Cf.  $\frac{F_{\mu\nu} F_D^{\nu\alpha} F_{\alpha}^{\mu}}{M^2}$
- Note that it is also naturally small, since it is induced by a loop. This whole story is special to U(1).

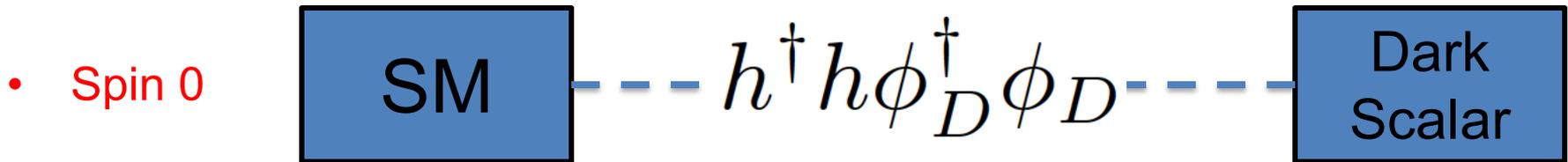
Okun (1982), Galison, Manohar (1984), Holdom (1986)

# DARK PHOTON, DARK HIGGS, STERILE NUS

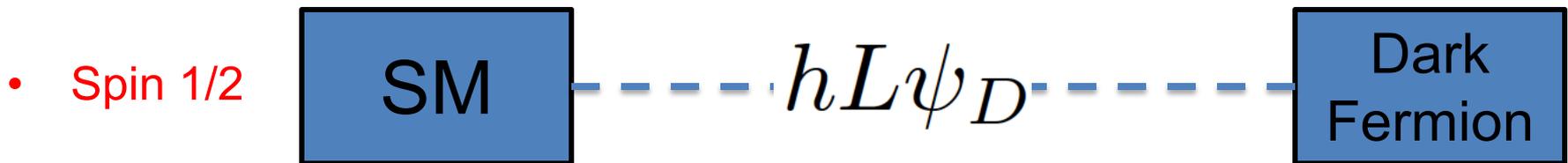
- This provides an organizing principle that motivates specific examples of new, weakly interacting light particles. There are just a few options:



→ **dark photon**, couples to SM fermions with suppressed couplings proportional to charge:  $\epsilon q_f$ . Holdom (1986)



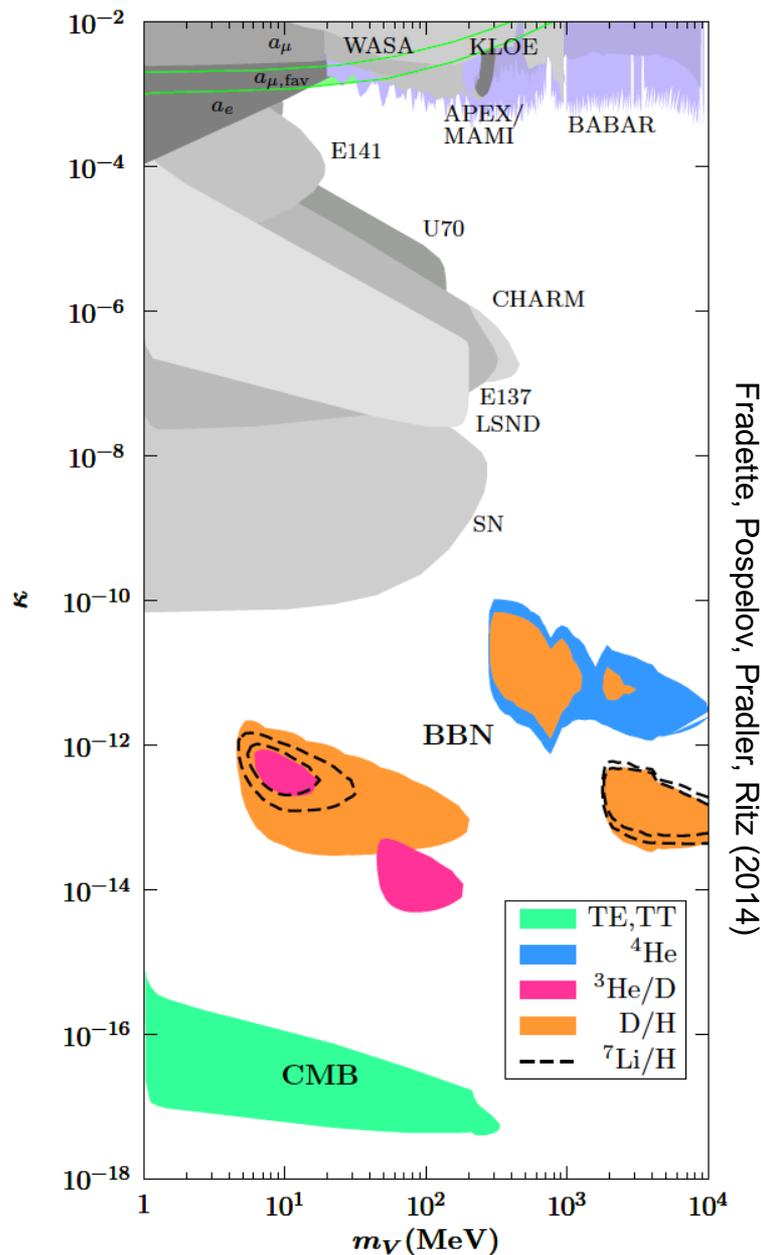
→ **dark Higgs boson**, couples to SM fermions with suppressed coupling proportional to mass:  $\sin \theta m_f$ . Patt, Wilczek (2006)



→ **sterile neutrino**, mixes with SM neutrinos with suppressed mixing  $\sin \theta$ .

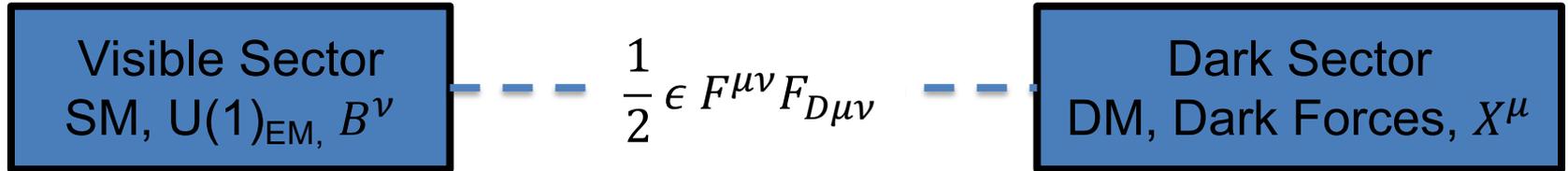
# EXISTING CONSTRAINTS

- We now have a few portal models, characterized by just a few parameters, and we can start exploring the parameter space.
- Consider dark photons. There is a vast parameter space.
  - “Bump hunts” exclude  $\epsilon > 10^{-3}$ .
  - Fixed target experiments exclude most of the gray region.
  - Astrophysics (supernova, BBN, CMB) excludes patches at very low coupling.
- But overall, light, weakly-interacting particles are much less constrained than  $\sim$ TeV, strongly-interacting particles.



# DARK PHOTON MODELS

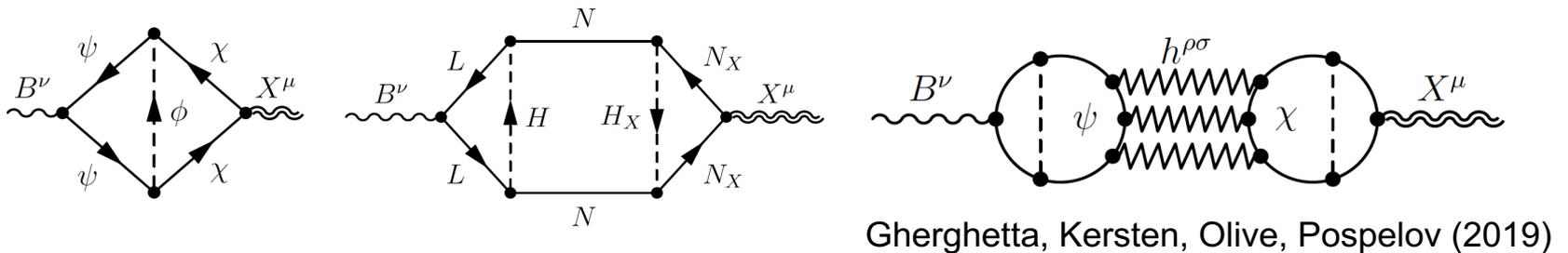
- If the dark photon is a portal particle, coupling arises from kinetic mixing:



- Mixing can be generated at 1-loop. If 0 at high scale, expect  $\epsilon \sim 10^{-3}$ .

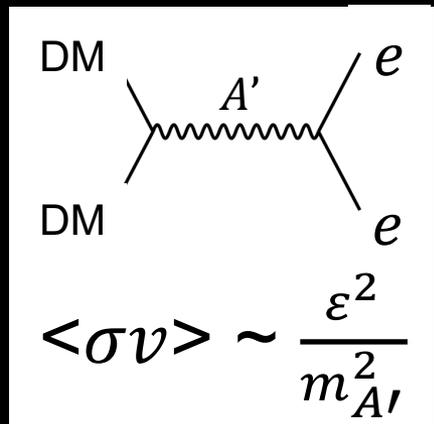
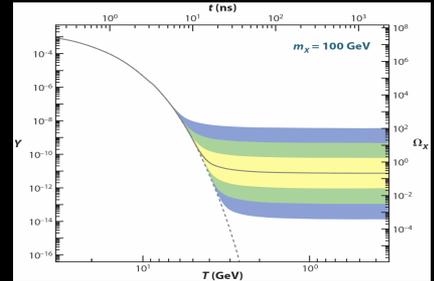
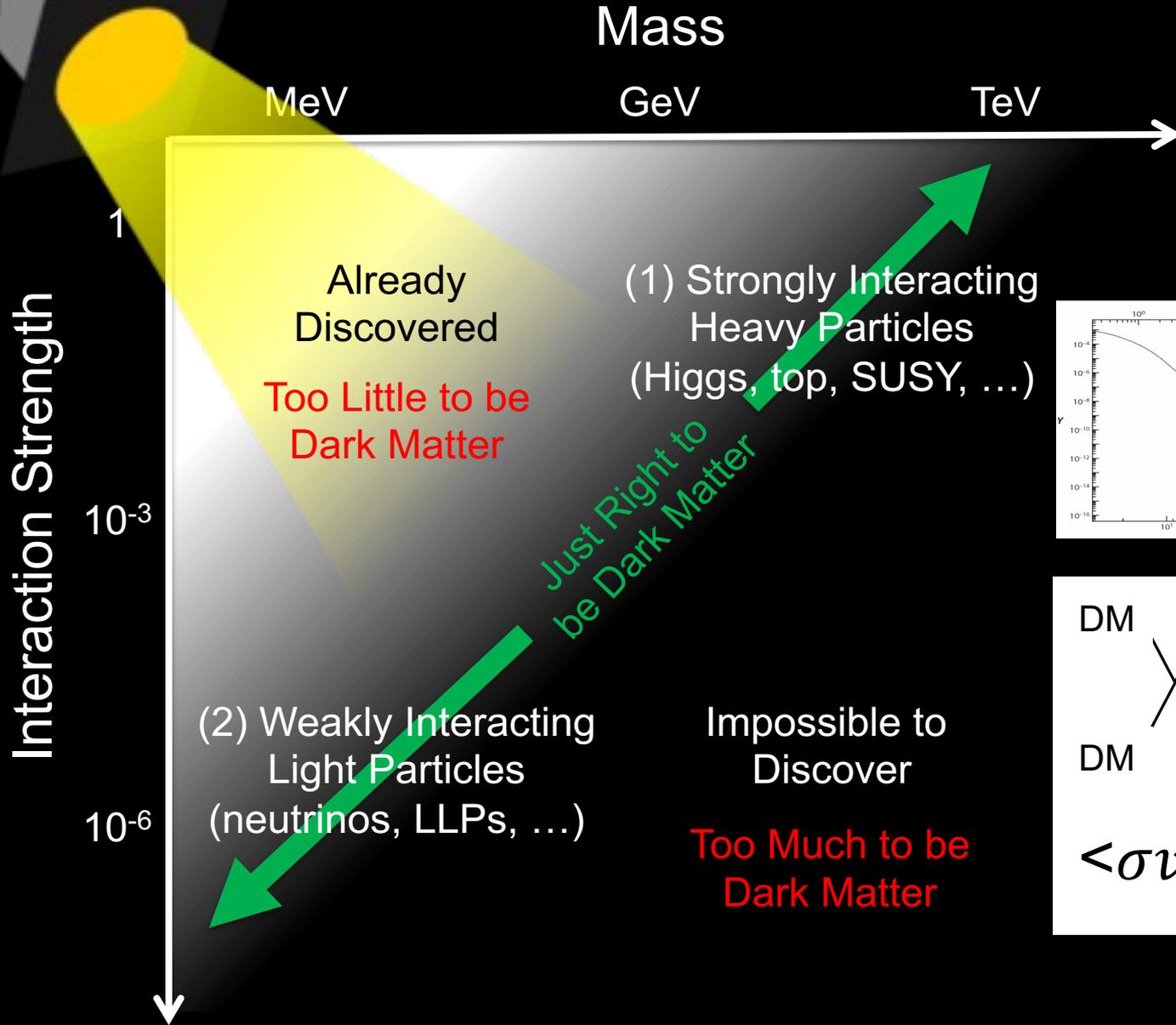
$$\epsilon = -\frac{g' g_X}{16\pi^2} \sum_i Y_i q_i \ln \frac{M_i^2}{\mu^2} \quad \text{Holdom (1986)}$$

- But there are also theories with mixing generated only at higher loop level.

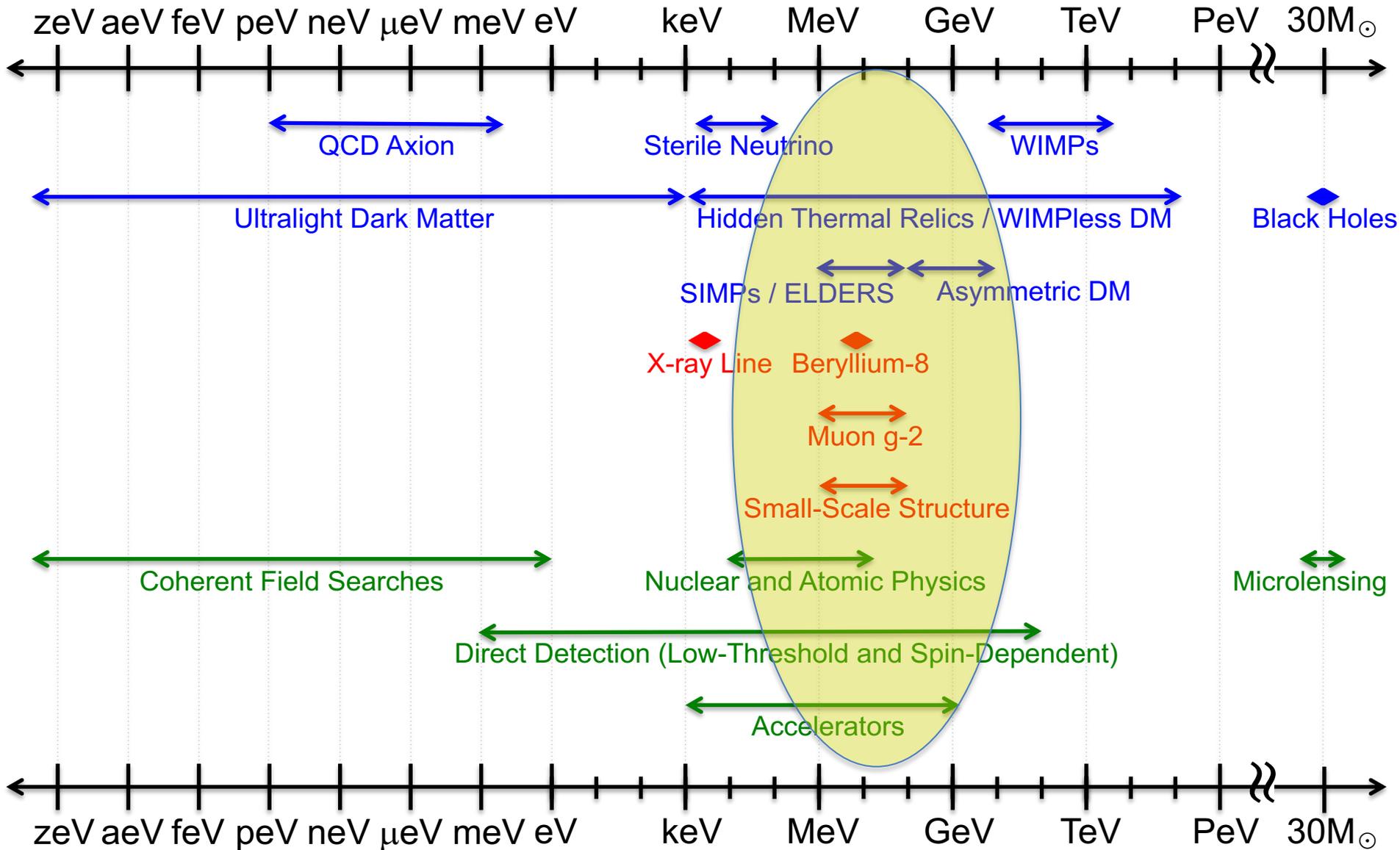


- Other than making us feel ok that  $\epsilon > 10^{-3}$  is excluded, models don't provide much guidance about the coupling, and none at all about the mass.

# DARK MATTER



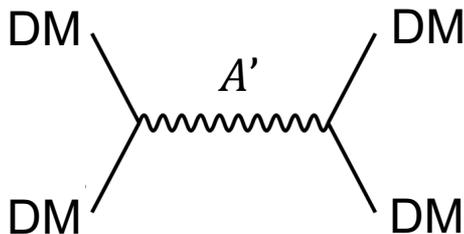
# Dark Sector Candidates, Anomalies, and Search Techniques



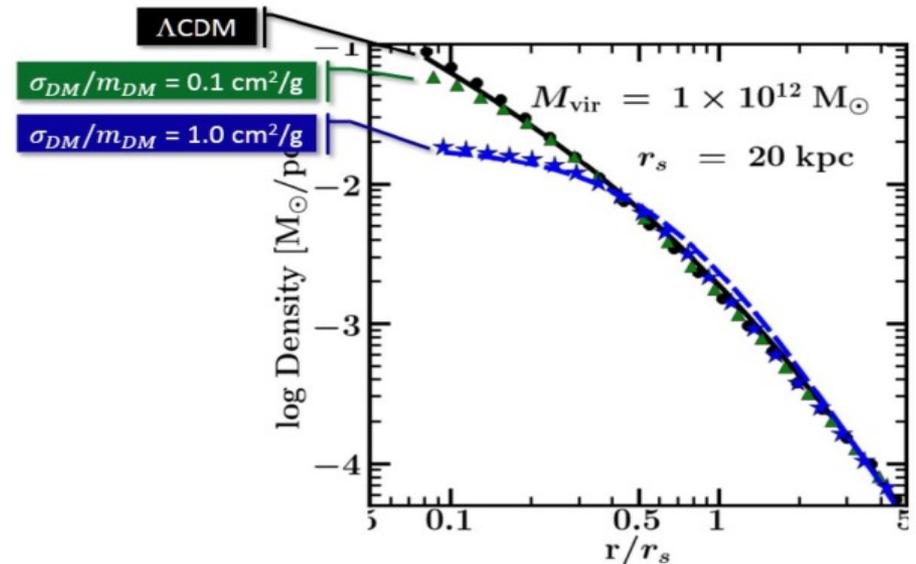
# SELF-INTERACTING DARK MATTER

- WIMP DM is in the “strongly interacting heavy particles” category. But there are indications that DM may not be WIMPs. No discovery so far, and also evidence from small-scale structure that dark matter may be strongly self-interacting.
- For example, there appear to be halo profiles that are not as cuspy (high central density) as predicted for standard collisionless cold dark matter (WIMPs, axions, sterile neutrinos, ...).
- To smooth out the cusps, need a self-interaction cross section

$$\frac{\sigma}{m} \sim \frac{\text{cm}^2}{\text{g}} \sim \frac{\text{barn}}{\text{GeV}} \sim (100 \text{ MeV})^{-3}$$



- This can be explained by a dark sector mass scale of  $\sim 10\text{-}100 \text{ MeV}$  (“dark neutrons interacting through dark pions”).

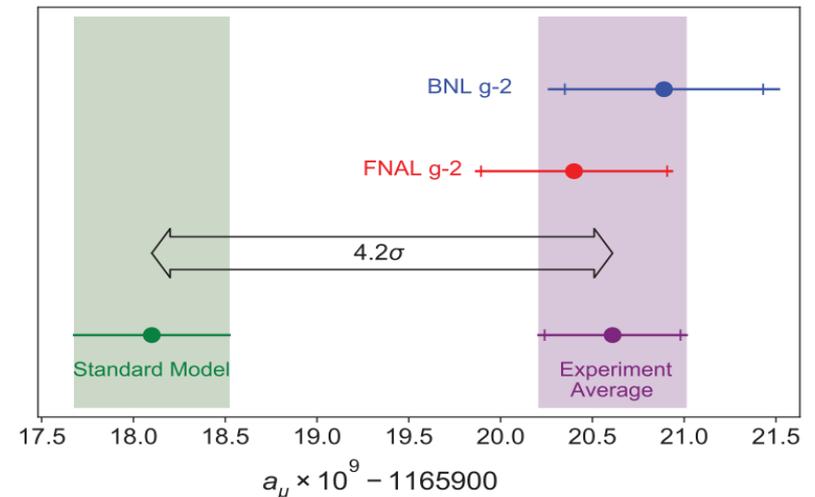
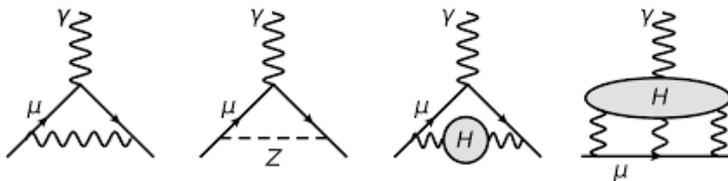
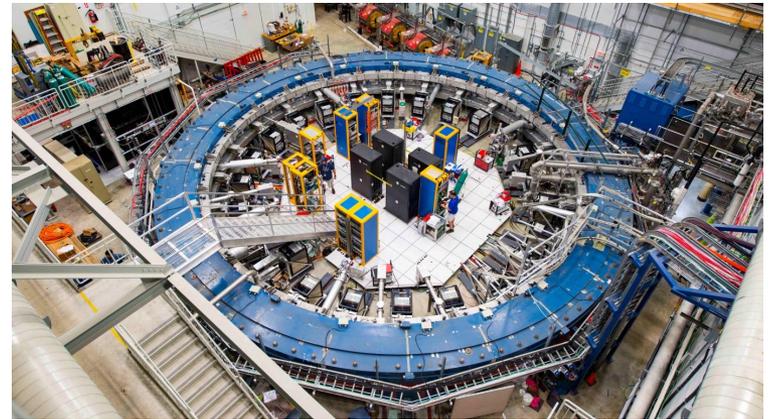


Tulin, Yu (2017)  
 Rocha et al. (2012); Peter et al. (2012)  
 Vogelsberger et al. (2012); Zavala et al. (2012)

# THE MUON'S ANOMALOUS MAGNETIC MOMENT

- In 2021, the Muon g-2 Collaboration announced a high precision measurement that deviates from the SM prediction by  $3.3\sigma$ .
- It is sensitive to the weak interactions, but unlike other precision probes, it requires neither flavor nor CP violation, and so is a “natural” place for new particles to appear, provided they couple to muons.

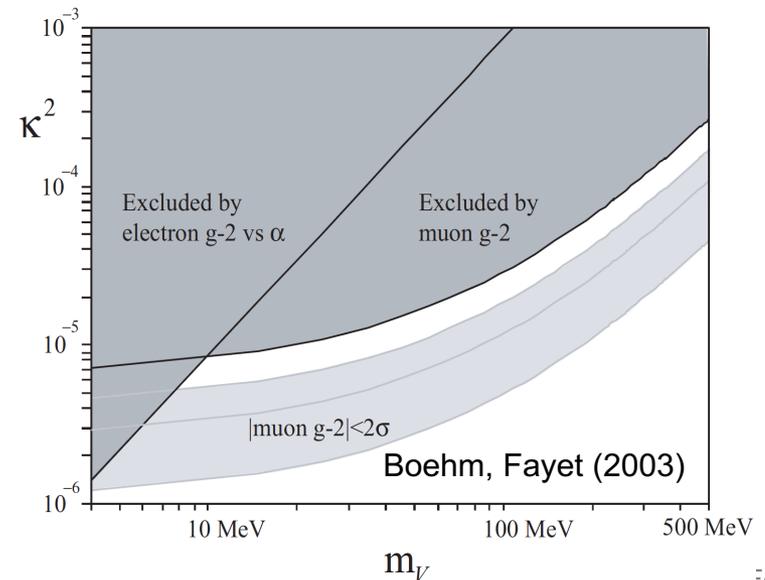
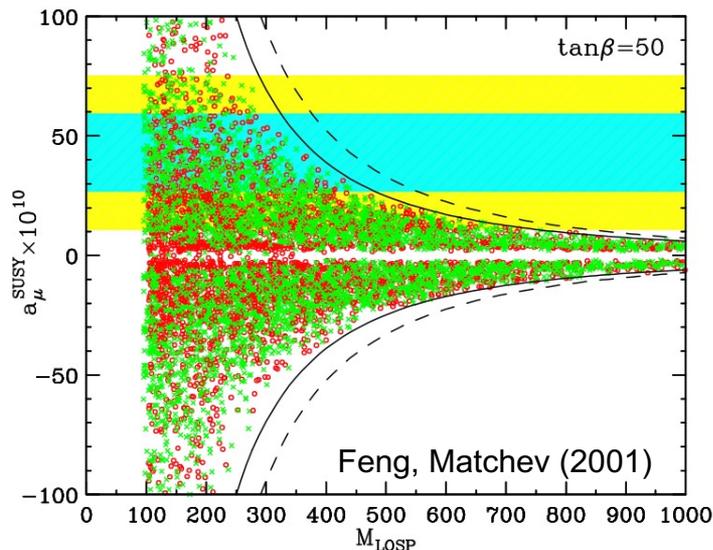
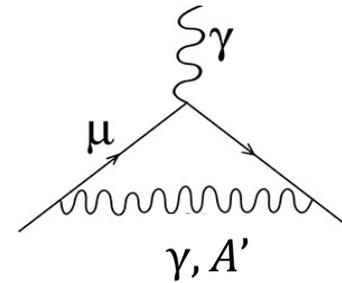
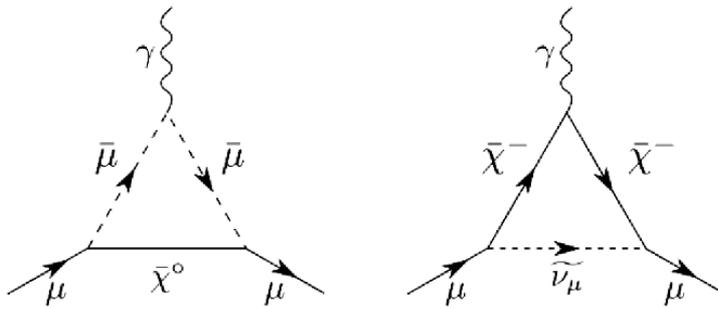
$$\vec{\mu} = g \frac{q}{2m} \vec{S} \quad a_l = (g_l - 2)/2$$



# THE MUON'S ANOMALOUS MAGNETIC MOMENT

- The discrepancy can be resolved by heavy particles, e.g., SUSY with superpartners at the 100s of GeV to TeV scale.

- But it can also be resolved with MeV-GeV masses and couplings  $\sim 10^{-3}$ . (Dark photon now excluded, but other similar particles remain viable.)



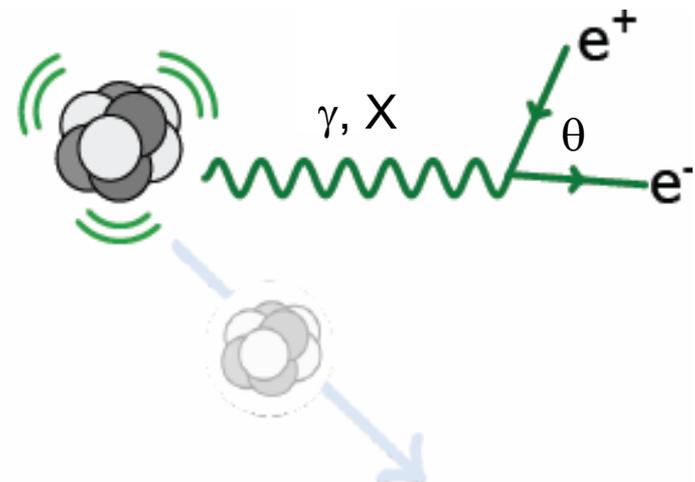
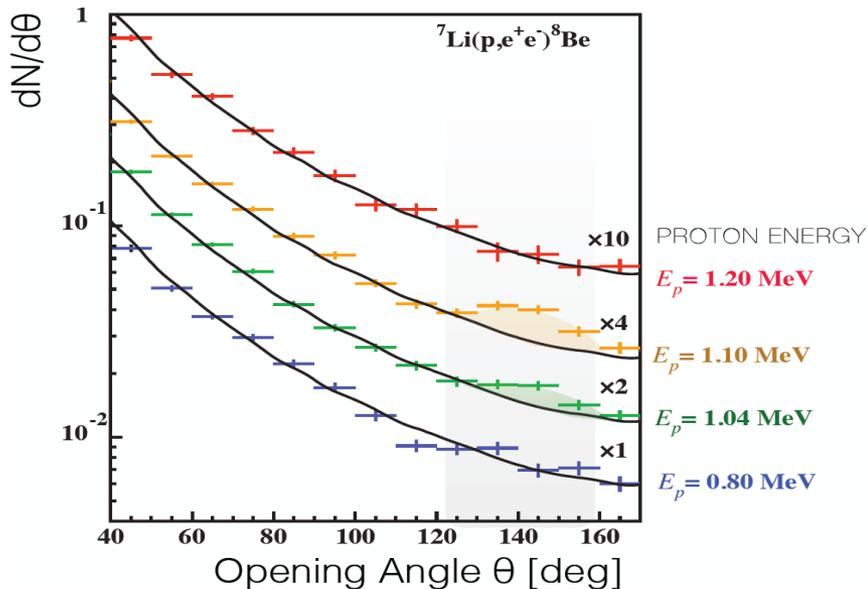
# THE ${}^8\text{Be}$ AND ${}^4\text{He}$ ATOMKI ANOMALIES

- New particles at the  $\sim 10$  MeV scale and below can be produced in the decays of excited nuclei.

Treiman, Wilczek (1978); Donnelly, Freedman, Lytel, Peccei, Schwartz (1978); Savage, McKeown, Filippone, Mitchell (1986)

- In 2015, an ATOMKI group reported a  $7\sigma$  excess in  ${}^8\text{Be}$  ( $18.15$ )  $\rightarrow {}^8\text{Be} e^+e^-$  decays at  $\theta_{e^+e^-} \approx 140^\circ$ .

Krasznahorkay et al., PRL, 1504.01527 [nucl-ex]



# THE ${}^8\text{Be}$ AND ${}^4\text{He}$ ATOMKI ANOMALIES

- The anomaly in the decays of excited  ${}^8\text{Be}$  nuclei can be explained by a new protophobic gauge boson  $X$  with mass 17 MeV and couplings  $\sim 10^{-4}$  to  $10^{-3}$ :  ${}^8\text{Be} (18.15) \rightarrow {}^8\text{Be} X$ , followed by  $X \rightarrow e^+ e^-$ .

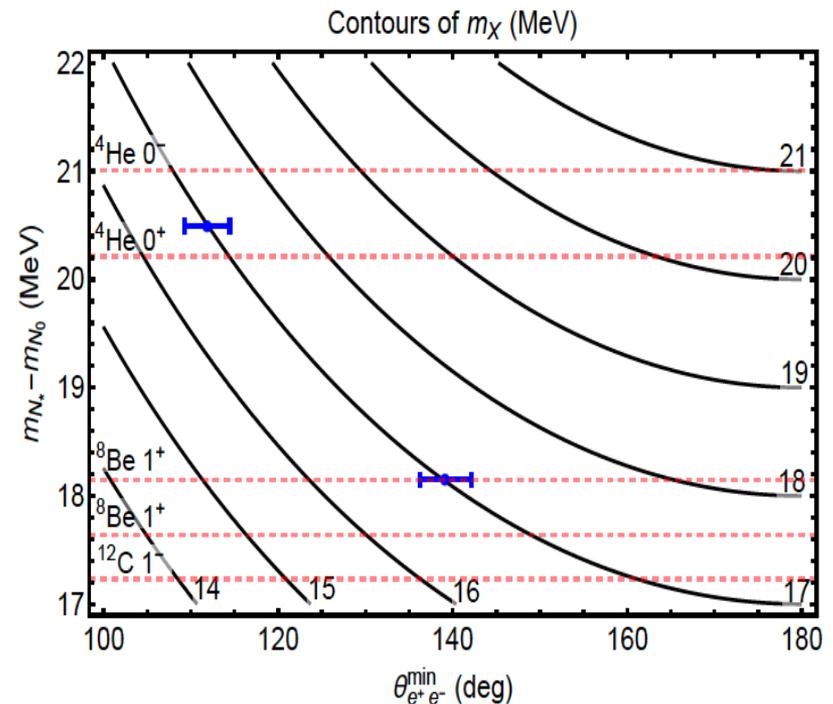
Feng, Fornal, Galon, Gardner, Smolinsky, Tanedo, Tait (2016)

- In 2019 the ATOMKI group reported a new  $7\sigma$  excess in the decays of excited  ${}^4\text{He} (20.49)$  nuclei at  $\theta_{e^+e^-} \approx 115^\circ$ .

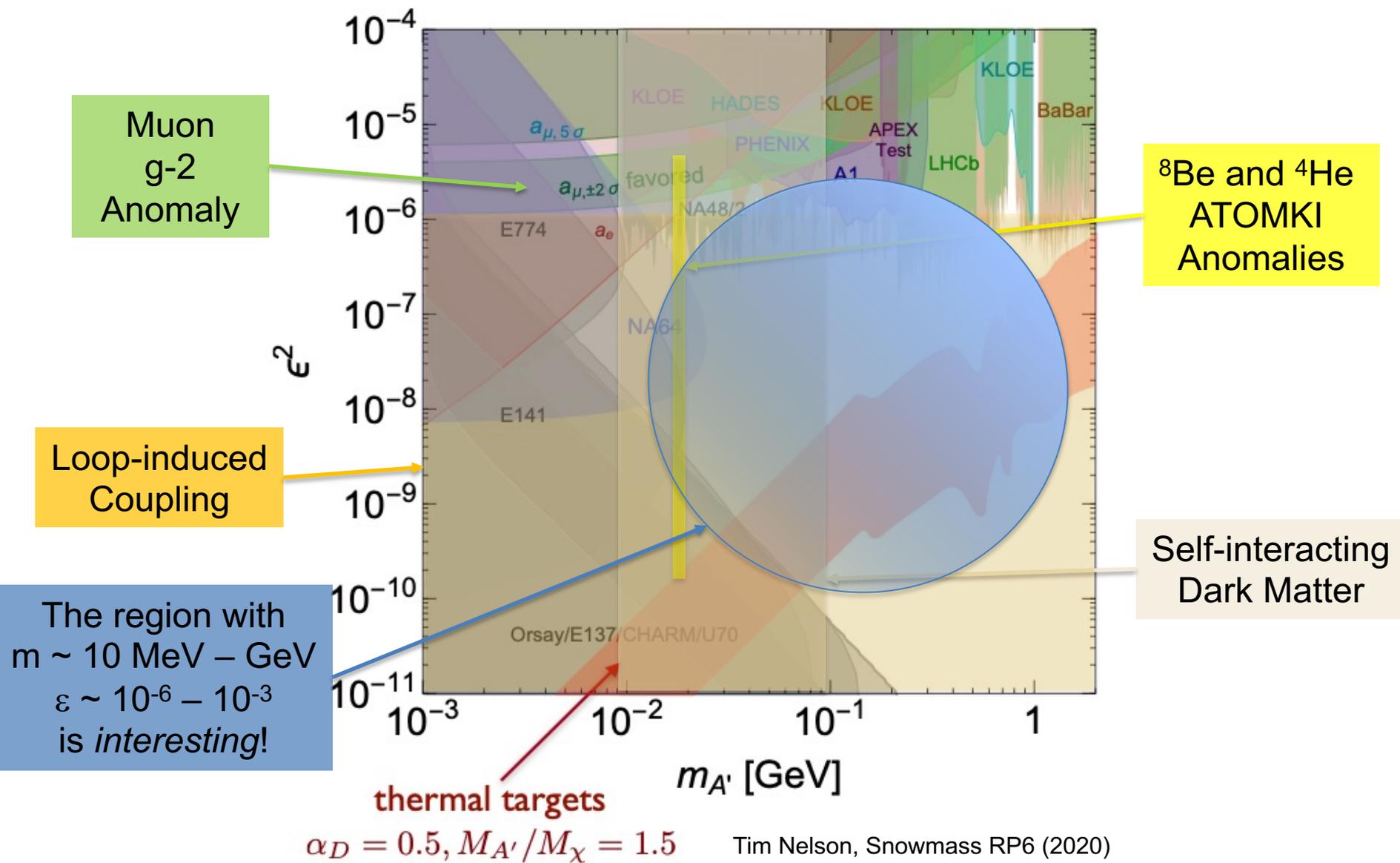
Krasznahorkay et al. (2019)

- Remarkably, this anomaly can be explained by the same new particle, which can also reduce the muon  $g-2$  discrepancy to  $2\sigma$ .

Feng, Tait, Verhaaren (2020)  
See also Zhang, Miller (2020)

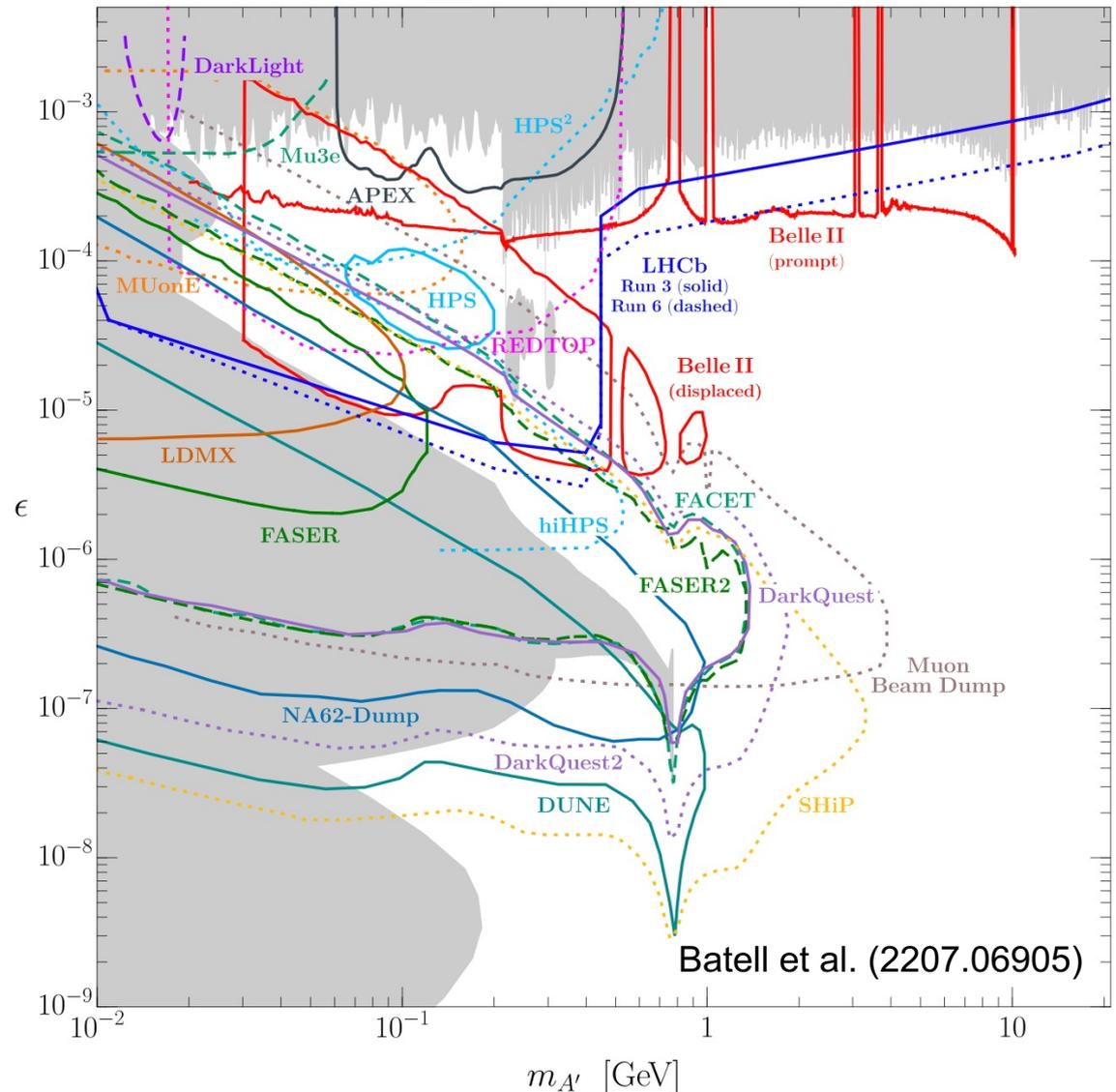


# TARGETS IN DARK PHOTON PARAMETER SPACE



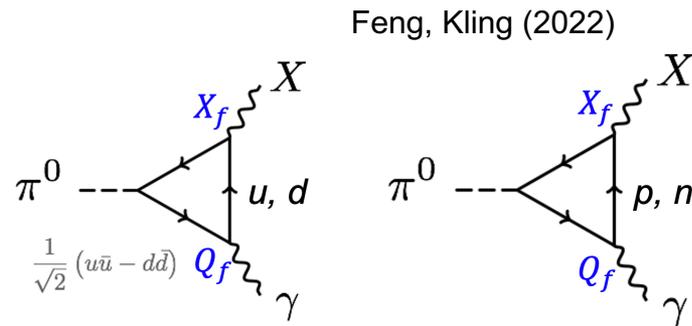
# CURRENT EXPERIMENTAL SEARCHES

- In the next few years, this region will be probed by currently running experiments (LHCb, Belle2, NA64, FASER, ...) and also proposed experiments. This is the low-hanging fruit of dark sectors – similar to Z-mediated WIMP cross sections. Soon this parameter space will look completely different.



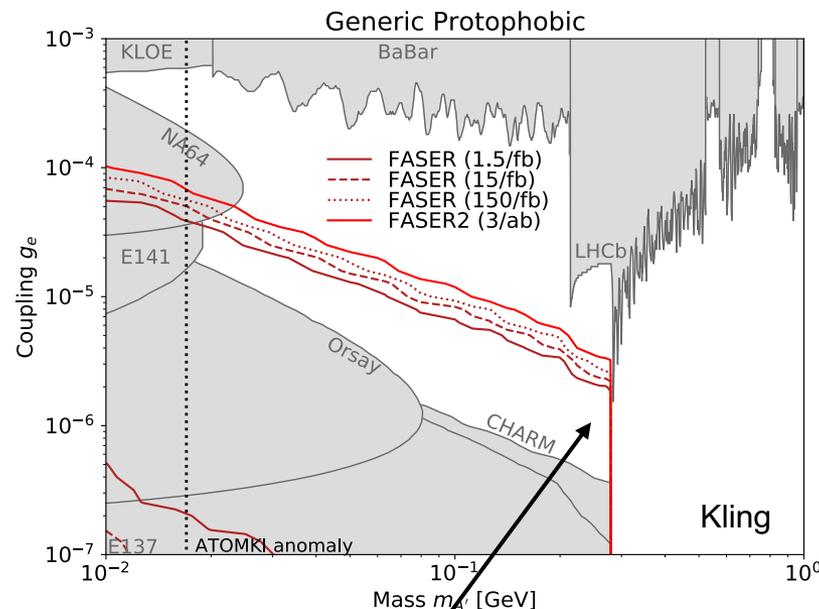
# PROTOPHOBIC GAUGE BOSONS

- Photophobic gauge bosons  $X$  are quite different from dark photons
  - No production by dark bremsstrahlung off protons
  - No production in pion decays: protophobic  $\rightarrow$  pion-phobic
  - Dominantly produced in  $\eta / \eta'$  decays.



This is eliminated if  $Q_u X_u - Q_d X_d \approx 0$  or  $2X_u + X_d \approx 0$  or  $X_p \approx 0$ .

- Consider a model-independent analysis: fix  $g_u, g_d \sim 10^{-3}$  (very large!) to explain ATOMKI, scan over  $g_e$ .
- With  $1 \text{ fb}^{-1}$ , FASER will probe the low  $g_e$  region. NA64 will probe high  $g_e$ .
- The 7-year-old ATOMKI anomalies will likely be first confirmed or refuted by high-energy experiments.

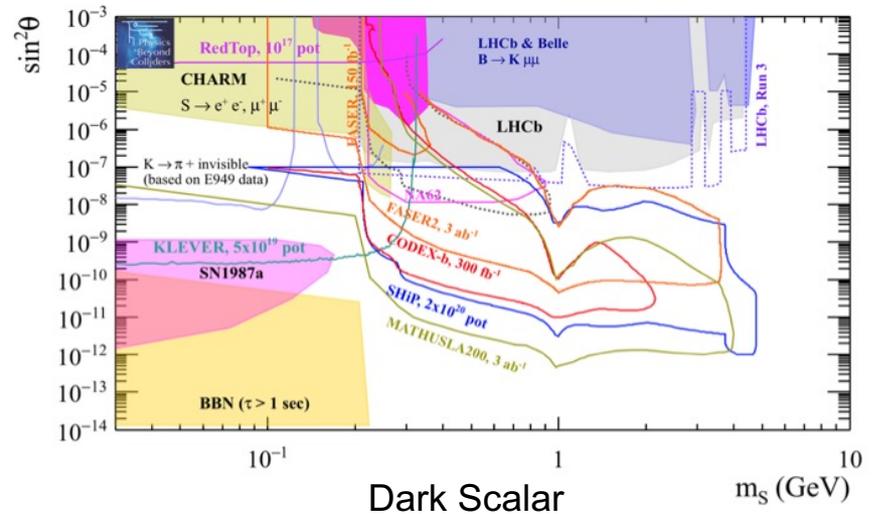
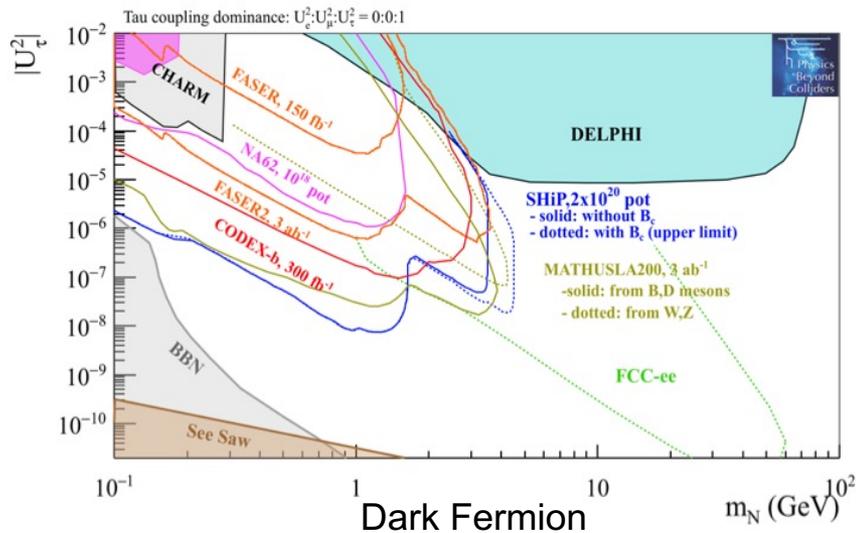
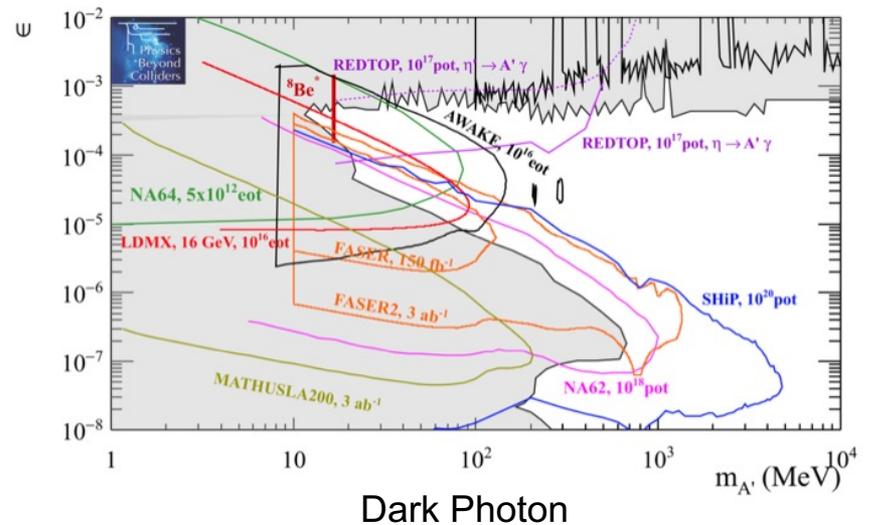


$$X \rightarrow \mu^+ \mu^- \text{ for } 2m_\mu < m_X < 2m_\pi$$

# FUTURE EXPERIMENTAL SEARCHES

- For the future, dedicated detectors have significant discovery potential for a wide variety of dark sector particles: dark photons; B-L and related gauge bosons; dark Higgs bosons; HNLs with couplings to e, mu, tau; ALPs with photon, gluon, fermion couplings; light neutralinos, inflatons, relaxions, and many others.

FPF White Paper (2022)

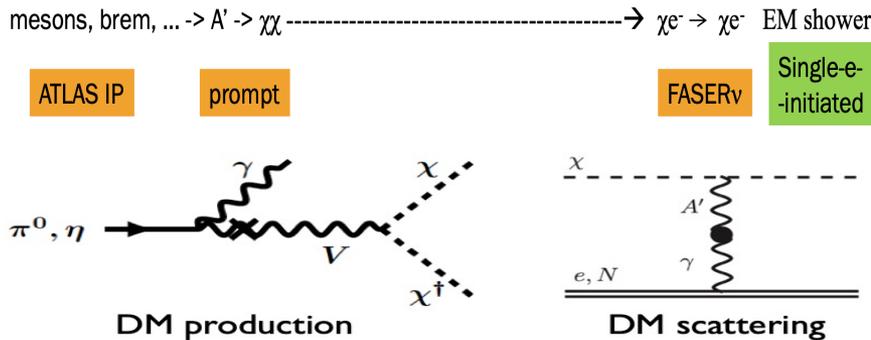


# DARK MATTER DIRECT DETECTION

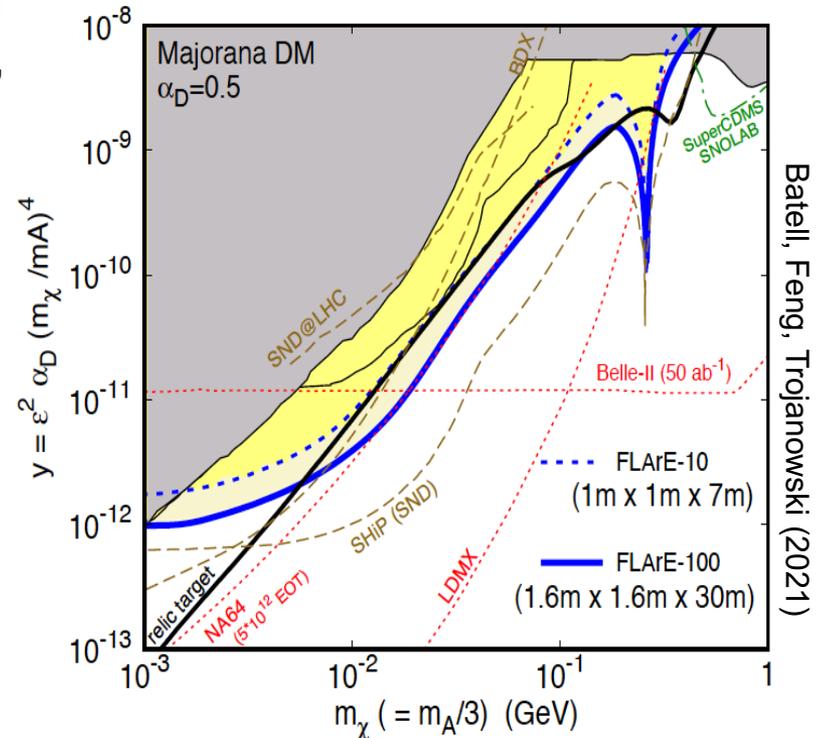
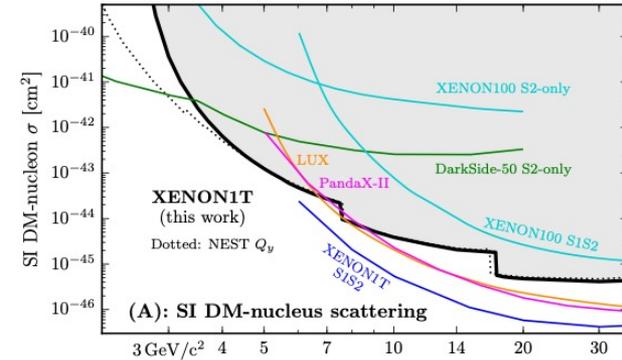
- What if the portal particle decays back to the dark sector? Light DM with masses at the GeV scale and below is famously hard to detect, but there is a great deal of creative work going on in this area.

See Asher Berlin's talk

- At the LHC, we can produce DM at high energies, look for the resulting DM to scatter in FLArE, Forward Liquid Argon Experiment, a proposed 10 to 100 tonne LArTPC.



- FLArE is powerful in the region favored/allowed by thermal freezeout. Note complementarity:  $q$  vs.  $e$  coupling, direct detection vs. missing  $X$ , etc.



# SUMMARY

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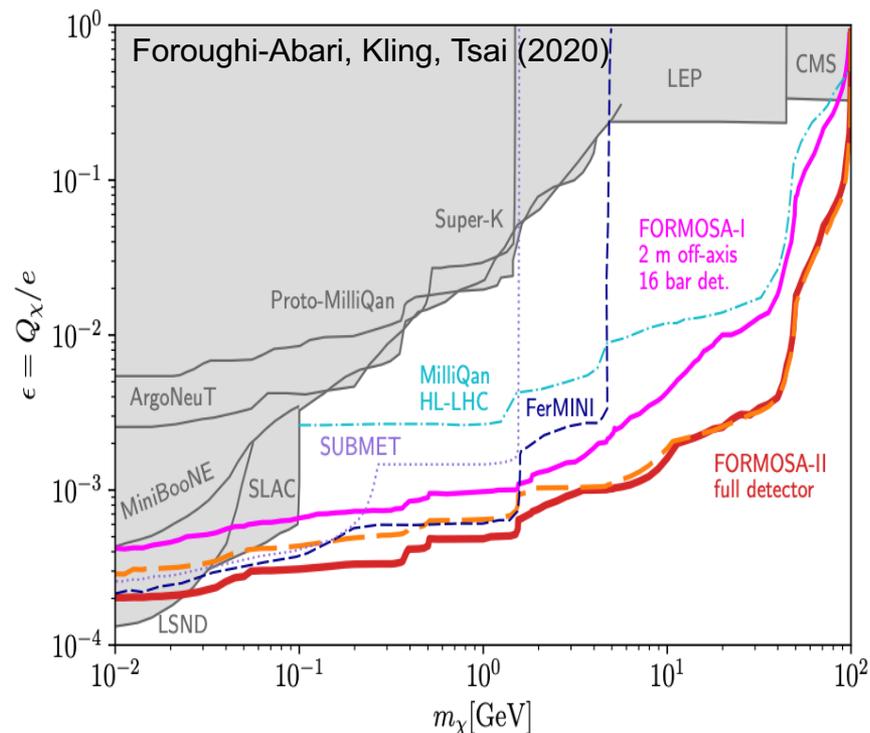
- Dark matter and dark sectors are currently among the leading motivations for BSM searches.
- In general, dark sectors predict new particles with an enormous range of masses and interaction strengths and also many qualitatively different possibilities; milli-charged particles, hadrophilic gauge bosons, quirks, ...; see backup slides.
- But there is a relatively small region of parameter space that has motivations comparable to WIMPs: masses  $\sim$  MeV to GeV, couplings  $\sim$  micro to milli.
  - WIMPs: coincidence of particle experiment (current threshold of what is observable), particle theory (gauge hierarchy problem), cosmology (WIMP miracle), astrophysics (large-scale structure), anomalies
  - Dark sectors: coincidence of particle experiment (current threshold of what is observable), particle theory (portals), cosmology (WIMPless miracle), astrophysics (small-scale structure), anomalies
- Most importantly, this region can be probed by many experiments in the coming 10 years. Just as WIMP DM became the subject of a world-wide research program in the 2000's and 2010's, dark sectors will become the subject of a world-wide research program in the 2020's and 2030's.

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# BACKUP

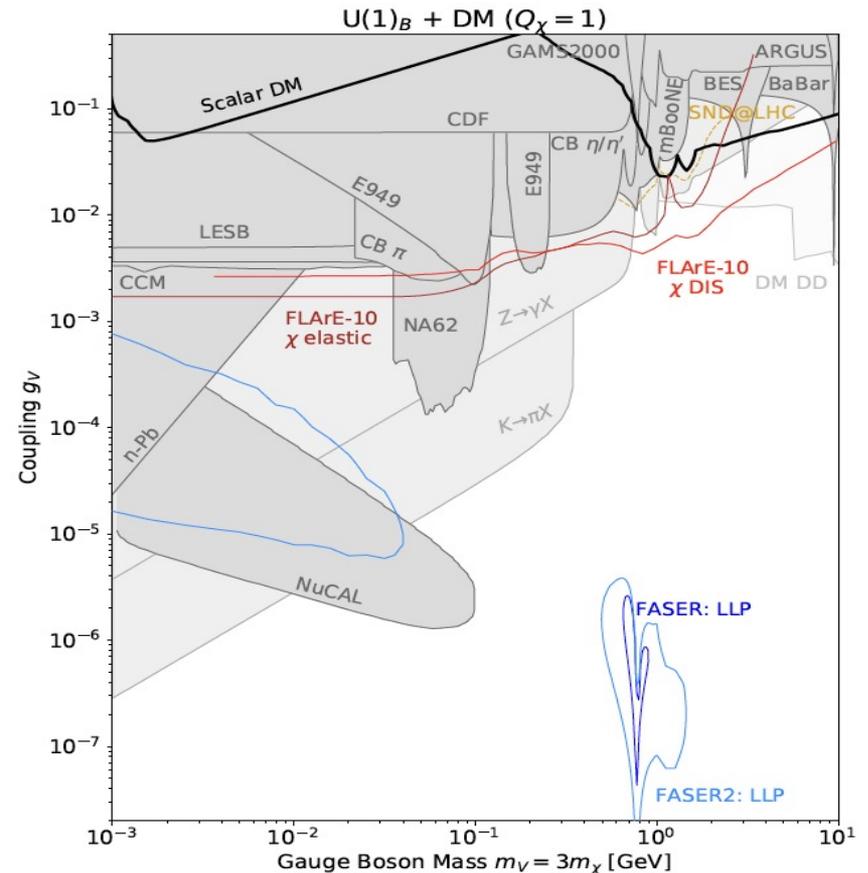
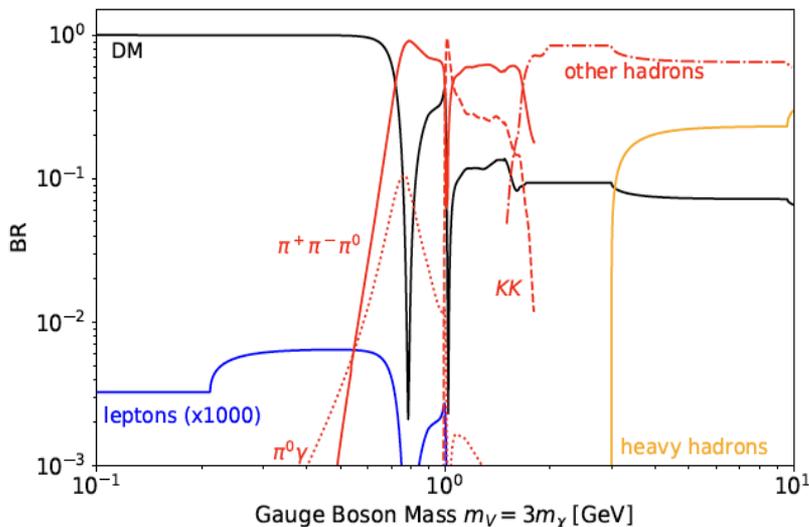
# MILLI-CHARGED PARTICLES

- A completely generic possibility motivated by dark matter, dark sectors. Currently the target of the MilliQan experiment, located at the LHC near the CMS experiment in a “non-forward” tunnel.
- The MilliQan Demonstrator (Proto-MilliQan) already probes new region. Full MilliQan can also run in this location in the HL-LHC era, but the sensitivity may be improved significantly by moving it to the FPF (FORMOSA).



# NEW SIGNALS: B AND B-3 $\tau$ GAUGE BOSONS

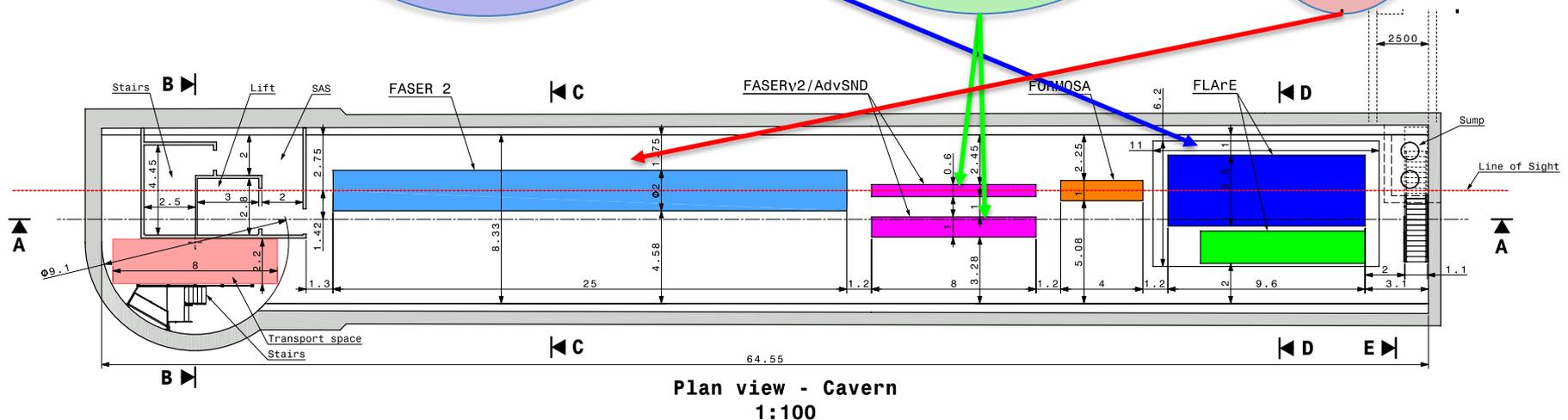
- Consider a light gauge boson coupled to baryon number
- Produced through  $q\bar{q} \rightarrow V$
- Many interesting hadronic decays  
 $V \rightarrow \pi^0\gamma, \pi^+\pi^-\pi^0, K^+K^-, K_S K_L$
- Greatly expands the standard  $e^+e^-$ ,  $\gamma\gamma$  signatures; similar signatures for “anomaly-free” gauge bosons



Batell, Feng, Fieg, Ismail, Kling, Abraham, Trojanowski, 2111.10343; see also Boyarsky, Mikulenka, Ovchinnikov, Shchutska, 2104.09688

# SIGNATURES FOR OTHER FPF EXPERIMENTS

Signature	DM DIS	DM Elastic	$\nu$ NC DIS	$\nu_\tau$ CC DIS	LLP decays
Models	$U(1)_B, U(1)_{B-3\tau}$	$U(1)_B, U(1)_{B-3\tau}$	$U(1)_{B-3\tau}$	$U(1)_{B-3\tau}$	$U(1)_B, U(1)_{B-3\tau}$
Production	$pp \rightarrow V \rightarrow \chi\chi$	$pp \rightarrow V \rightarrow \chi\chi$	$pp \rightarrow D_s \rightarrow \nu_\tau$	$pp \rightarrow V \rightarrow \nu_\tau \bar{\nu}_\tau$	$pp \rightarrow V$
Detection	$\chi N \rightarrow \chi X$	$\chi P \rightarrow \chi p$	$\nu_\tau N \rightarrow \nu_\tau X$	$\nu_\tau N \rightarrow \tau X$	$V \rightarrow \text{hadrons}$



# NEW SIGNALS: QUIRKS

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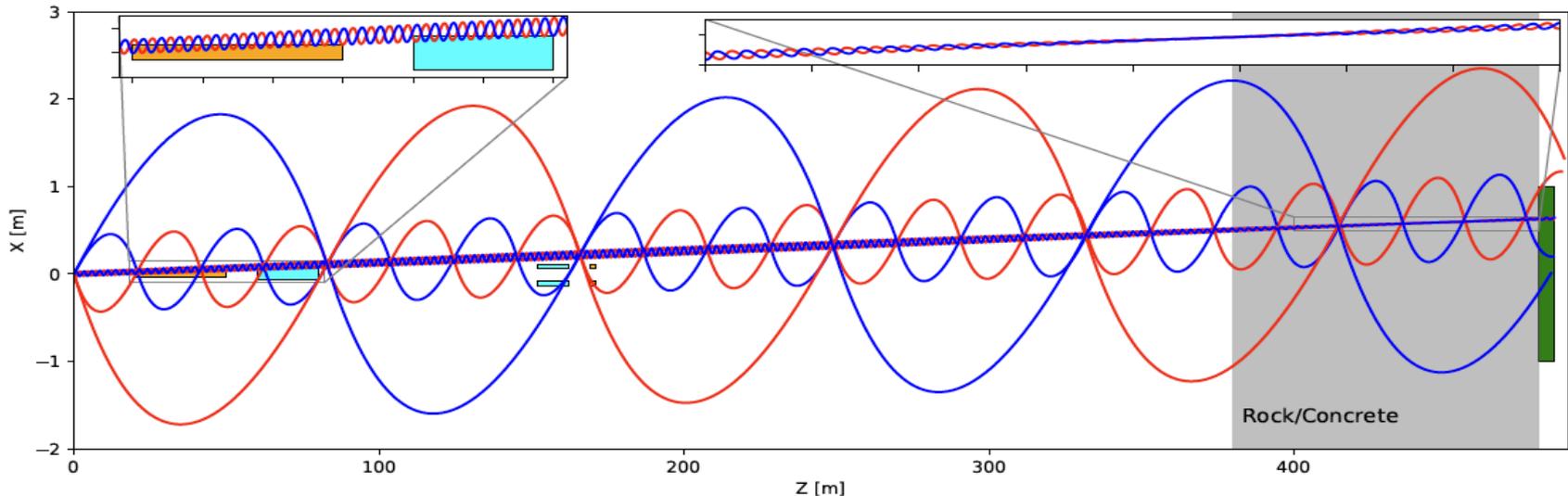
Kang, Luty (2008)

- Quirks are matter particles charged under a hidden strong force with mass  $m \gg \Lambda_{\text{hidden}}$ . E.g.,  $m \sim 100 \text{ GeV} - \text{TeV}$ ,  $\Lambda_{\text{hidden}} \sim \text{keV}$ .
- Quirks may also have SM charge and color. They are then pair produced at the LHC, and are connected by a hidden color string.
- For quarks and standard QCD,  $m \ll \Lambda_{\text{QCD}}$ , and so it becomes energetically favorable to pair produce new quarks from the vacuum. Quarks hadronized.
- But for quirks, since  $m \gg \Lambda_{\text{hidden}}$ , it is never energetically favorable to break the string by pair producing quirks from the vacuum: quirks do not hadronize, they oscillate.

# QUIRK SIGNATURE

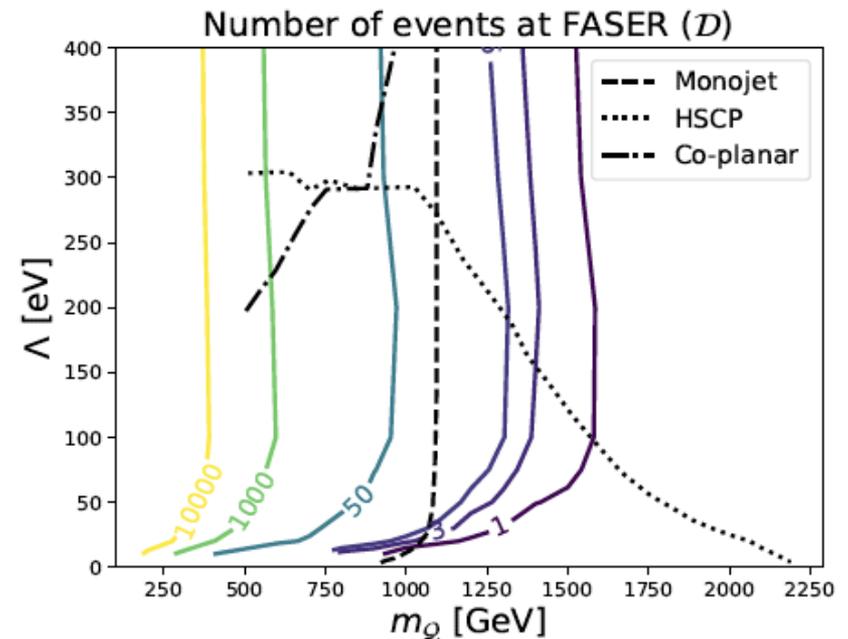
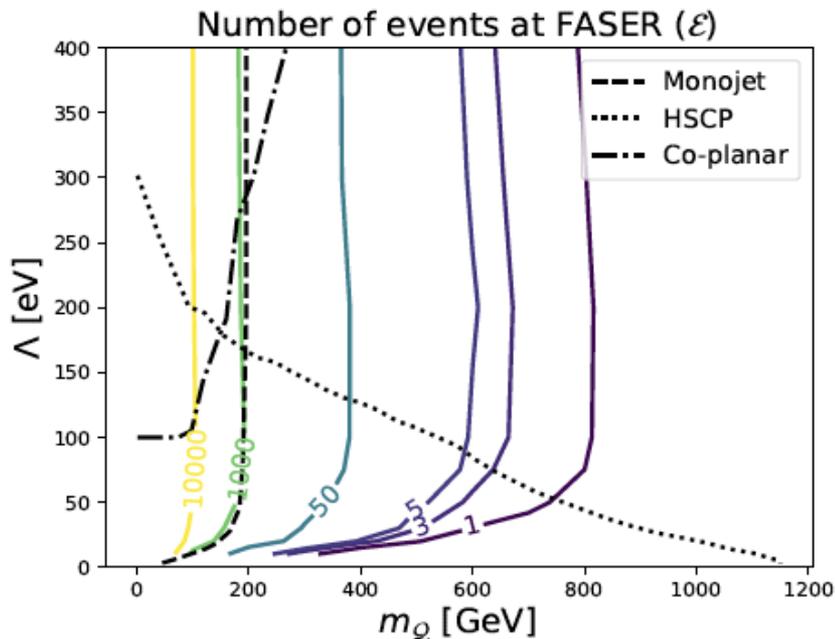
- Of course, the quirk – anti-quirk system has low  $p_T$ .
- The pair therefore oscillates, with length scale  $\sim 1/\Lambda_{\text{hidden}}$ .
- For a range of  $\Lambda_{\text{hidden}}$ , the quirk system travels down the beamline, escaping most LHC detectors, but ultimately leaving (strange!) tracks in FASER.

Li, Pei, Ran, Zhang, 2108.06748



# QUIRK DISCOVERY PROSPECTS

- Far-forward detectors at the LHC are ideally suited to search for quirks.
  - Like heavy particles, they require the LHC to be produced
  - Like light particles, they are dominantly produced along the beamline
- $\sim 1000$  of events possible at FASER in Run 3



Li, Pei, Ran, Zhang, 2108.06748